

4 CP Quantum Mechanics

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4.1 Basic Assumptions

Modeling of CP is based on the following basic assumptions:

- (1) CP contain ensembles of atomic nuclei densely lined up in a long and very narrow channel.
- (2) The distances between the nuclei are so small, that all electrons bound to these nuclei are delocalized along the channel. In other words: Even in their electronic ground state CP don't consist of individual atoms. CP rather form a quasi-one-dimensional plasma (this could also be seen as a metal).

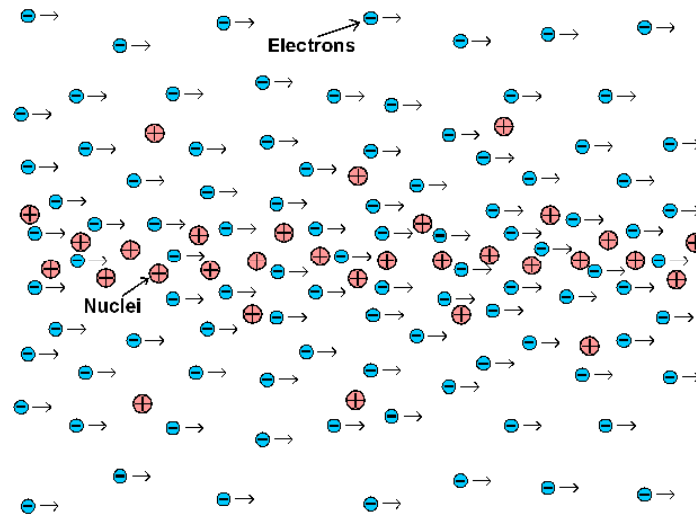


Figure 3 Basic model of a CP. The CP similarly extends to the left and to the right of this picture

4.2 The Cylindrical Model of CP

The shape and quantum mechanical state of CP can be very complicated. In order to obtain a simple quantum mechanical description of CP, the following simplifications are used, which will subsequently be called "the cylindrical model of CP":

- (3) The CP is perfectly straight and cylindrically symmetric, i.e. it is not bent to rings, helices etc. The CP is oriented in parallel to and centered on the z-axis of the modeling cylindrical coordinate system.
- (4) The CP has the length \bar{L} and contains a total nuclear charge Q .
- (5) The electron wave functions of the CP are confined in the interval $0 \leq z < \bar{L}$. At $z = \bar{L}$ these wave function are continuously extended to their value and gradient at $z = 0$, as if the CP were rings. However, this is meant to describe only the *circular boundary condition* of the wave functions at $z = \bar{L}$, not the shape of the CP.
- (6) No external field is applied to the CP.
- (7) A *jellium* model is used for the spatial distribution of the nuclear charge. This means, for the purpose of computing the spatial distribution of the electrons, the positive charges of the nuclei are modeled as a uniform "positive jelly" background, rather than point charges with distances in between. The nucleic charge density is assumed to be constant in axial and azimuthal direction, but it depends on the radial distance.
- (8) However, correction terms will be used for computing the long-range Coulomb interaction, which are taking care of the jellium's short-range granularity (i.e. its point charge quality).
- (9) The nucleic charge distribution of the jellium is modeled by means of a two-dimensional normal distribution in radial direction (or alternatively, by a distribution function, which is more aligned with the computed electron charge distribution). The standard deviation (or the alternative distribution function) is to be determined by variation, such that the total energy of the CP is minimized.

- (10) The core area of the CP (i.e. the space defined by the extend of its electron orbitals) is surrounded by a charge compensation layer (“halo”) consisting of cations. The charge of the halo compensates the surplus negative charges of the electrons in the core. The halo is modeled as a cylindrical shell of positive charge. The halo radius has to be larger than the core radius.
- (11) The CP is assumed to reside in a vacuum. Interaction of the CP with surrounding matter is thus neglected.
- (12) Only stationary states are modeled, as the goal is to describe the ground state of CP. Consequently, the model assumes there is no electron scattering, i.e. there is no momentum transfer between electrons and the nuclei.
- (13) For computing the repulsion energy among the nuclei, short-range corrections to the jellium model have to be made, which account for the granularity of the nuclear charges. In case the CP contains a mixture of different sorts of atomic nuclei, only the mean nuclear charge is taken into account for the corrections, rather than the individual nuclear charges.
- (14) The time-independent Klein-Gordon equation is used for modeling the electron wave functions, thereby neglecting the magnetic moments of electron spins. The Klein-Gordon equation is taking care of the large relativistic effects occurring in CP. (Clearly, the Dirac equation would be more adequate for modeling CP. However, the involved complexities of such approach are avoided here.). For comparing the formulas and simulation results with the ones obtained from a non-relativistic Hamiltonian, also the Schrödinger equation is used.
- (15) The magnetic field of the azimuthal electron orbits is neglected.
- (16) Magnetic field from nuclear spins is neglected.
- (17) The electron wave functions are modeled in an inertial frame of reference, where no magnetic field is created by any collinear movements of the nuclei. This simplification amounts to an approximation in cases where the nucleic velocities are position dependent.
- (18) The multi-electron system is approximated by computing a collection of one-electron orbitals, whereby each electron orbital is subjected to the mean electric potential and magnetic vector potential created by the total charge density and total current density of all other occupied orbitals and the nuclei (independent particle model). The Pauli exclusion principle is used for determining orbital occupations of the ground state. Exchange and correlation energies are neglected.
- (19) Quantum field theory is not engaged. Particle count is conserved.
- (20) Eigenstates are excluded from occupation, where the corresponding total energy eigenvalue (including the electron’s rest energy) of the electron is negative. This shall ensure that the mass defect per electron doesn’t exceed the electrons rest energy.
- (21) The kinetic energy of the electrons is always positive, i.e. states with a negative kinetic energy are ignored.
- (22) Only bound eigenstates of the electrons are considered, i.e. the total energy of an eigenstate has to be less than the electron rest energy (i.e. the sum of the potential energy and the kinetic energy has to be negative).

4.3 The Klein-Gordon Equation of a CP

Initial calculations of a CP with the Schrödinger equation have shown, that the resulting binding energy of the electrons would potentially far exceed the rest energy of the electrons. Also, the axial electron velocities would potentially far exceed the speed of light. These results were absurd in light of special relativity, because the mass defect per bound electron should never exceed twice its rest mass and velocities should always stay lower than the speed of light. This provided ample reasons for engaging a relativistic Hamiltonian and a Lorentz-covariant quantum mechanical equation to model CP.

Generally, the Dirac equation is regarded as the correct Lorentz-covariant equation for modeling fermions, especially when the effects resulting from the particle’s spin is of concern. Unfortunately, the Dirac equation involves 4-component wave functions and the solution of four coupled differential equations, resulting in sizeable mathematical and computational efforts.

Assuming that the electron spins have only minor effects on the binding energy, charge density, current density and other observables, the Klein-Gordon equation provides a Lorentz-covariant alternative to the Dirac equation for modeling the electrons of CP. At the non-relativistic limit the Klein-Gordon equation is equivalent to the Schrödinger equation, while both equations share the deficiency of not modeling the spin.

In relativistic electrodynamics with so-called minimal coupling the Hamiltonian (total energy) of a particle with charge q moving in the presence of a static (external) electromagnetic potential is:

$$(23) \quad \hat{H} = c\sqrt{(\vec{P} - q\vec{A})^2 + m_e^2 c^2} + q\Phi, \text{ where } c \text{ is the speed of light, } \Phi \text{ is the electric potential, } \vec{A} \text{ is the magnetic vector potential, } \vec{P} = \gamma m_e \vec{v} + q\vec{A} \text{ is the electron's canonical momentum, } m_e \text{ is the electron rest mass and } \gamma \text{ is the Lorentz factor}$$

By defining $\bar{E} \equiv \hat{H} - m_e c^2$ as being the total energy minus the rest energy and by using $q = -e$ as the charge of an electron, (23) is leading to the following equation for an electron in a static electromagnetic potential:

$$(24) \quad \bar{E} = c\sqrt{(\vec{P} + e\vec{A})^2 + m_e^2 c^2} - m_e c^2 - e\Phi, \text{ where } e \text{ is the elementary charge}$$

Therefore:

$$(25) \quad (\bar{E} + e\Phi + m_e c^2)^2 = (c\vec{P} + ec\vec{A})^2 + m_e^2 c^4$$

All formulas are written in SI units, unless otherwise noted. Throughout this document, energy symbols with a bar on top (e.g. \bar{E}) denote, that the energy is measured in Joule. Energy symbols without a bar on top denote, that the energy is measured in units of the Hartree energy (117), i.e. the energy is a dimensionless quantity in the respective formula. Likewise, other symbols with a bar (e.g. $\bar{\sigma}$, \bar{J}_z , \bar{p}_z , \bar{P}_z , \bar{A}_z) are in SI units, while its counterparts without the bar are in natural units (i.e. dimensionless).

By quantizing the canonical momentum via the del operator $\vec{P} \equiv -i\hbar\nabla$ and applying both sides to an electron wave function Ψ , equation (25) transforms to the **stationary Klein-Gordon equation of an electron in a static electromagnetic potential**:

$$(26) \quad (\bar{E} + e\Phi + m_e c^2)^2 \Psi = \left[(-i\hbar c \nabla + ec\vec{A})^2 + m_e^2 c^4 \right] \Psi, \text{ where}$$

\hbar is the reduced Planck constant and $i = \sqrt{-1}$

Due to simplification (19), Ψ is called here a “wave function”, rather than a “quantum field”.

The term $\bar{E} + m_e c^2$ represents the total energy of the electron. Usually the Klein-Gordon equation is written, such that the total energy is sought as the eigenvalue of this differential equation. However, this document deviates from the customary approach. Instead, the quantity \bar{E} is sought here as the eigenvalue (both approaches are equivalent in their results).

Some authors prefer the term “relativistic Schrödinger equation” for (26), insisting that the Klein-Gordon equation is different. Here, these terms are used interchangeably.

In quantum mechanics a multi-electron system is correctly described by a single wave function $\Psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N)$ depending on the positions of the N electrons. The multi-electron wave function is usually formed by a Slater determinant (or a linear combination of several Slater determinants) to ensure anti-symmetry and the Pauli exclusion principle.

However, the number of electrons in a CP can exceed 10^{10} , which renders a Slater determinant entirely impractical to compute, because a program cannot handle equations with e.g. 10^{10} positions and compute determinants of this size.

According to simplification (18) a rigorously simpler approach is used here for modeling CP, requiring only moderate compute power:

So, instead of using a multi-electron Klein-Gordon equation describing the pair-wise interaction between N electrons, the cylindrical model uses N **single-electron Klein-Gordon equations with N wave functions** $\Psi(\vec{r})$, each describing a single electron in the **mean potential** of all other electrons and the nuclei.

Of course, this is merely an independent particle approximation. For example, the approach doesn't account for the exchange energy and the correlation energy usually deemed important in quantum chemistry.

At first glance this looks still challenging to compute, because there are N Klein-Gordon equations to be solved. Fortunately, large numbers of these equations can be computed in groups, because they produce nearly the same charge density distributions and current density distributions.

Expanding the right side of (26) and using $\nabla \cdot \vec{A} = 0$ (Lorentz gauge in the static case) yields:

$$(27) \quad (\bar{E} + e\Phi + m_e c^2)^2 \Psi = \left(-\hbar^2 c^2 \nabla^2 - 2i\hbar e c^2 \vec{A} \cdot \nabla + e^2 c^2 \vec{A} \cdot \vec{A} + m_e^2 c^4 \right) \Psi$$

A kinetic momentum operator is defined here as:

$$(28) \quad \hat{p} = -i\hbar \nabla + e\vec{A}$$

The expectation value of \hat{p} equals $\langle \gamma \rangle m_e \langle \vec{v} \rangle$, hence the name “kinetic momentum” ($\langle \vec{v} \rangle$ is the expectation value of the electron’s group velocity and $\langle \gamma \rangle = \left\langle 1 + (\bar{E} + e\Phi)/(m_e c^2) \right\rangle$ is the expectation value of the local Lorentz factor).

Using (28) in (27) yields:

$$(29) \quad (\bar{E} + e\Phi + m_e c^2)^2 \Psi = \left(c^2 \hat{p}^2 + m_e^2 c^4 \right) \Psi, \text{ where } c^2 \hat{p}^2 = -\hbar^2 c^2 \nabla^2 - 2i\hbar e c^2 \vec{A} \cdot \nabla + e^2 c^2 \vec{A} \cdot \vec{A}$$

According to simplification (14) and (15) the magnetic field of the electron spins and of the azimuthal movement of the electrons is neglected. Thus the only source of the magnetic field is the current carried by the electrons moving in z-direction. Therefore, the vector potential is everywhere oriented in z-direction:

$$(30) \quad \vec{A} = \bar{A}_z \vec{e}_z$$

The Laplace operator expands in cylindrical coordinates as following:

$$(31) \quad \nabla^2 = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2}, \text{ where } \rho \text{ is the radial distance from the z-axis, } \varphi \text{ is the azimuth and } z \text{ is the coordinate of the z-axis}$$

Inserting (30) and (31) into equation (27) and dividing both sides by $2m_e c^2$ is resulting in **the stationary Klein-Gordon equation of an electron in the mean potential of a CP’s all other electrons and the nuclei:**

$$(32) \quad \left\{ \frac{-\hbar^2}{2m_e} \left[\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2} + 2 \frac{e\bar{A}_z}{\hbar} i \frac{\partial}{\partial z} - \frac{e^2 \bar{A}_z^2}{\hbar^2} \right] - \frac{m_e c^2}{2} \left(\frac{\bar{E} + e\Phi}{m_e c^2} + 1 \right)^2 + \frac{m_e c^2}{2} \right\} \Psi = 0$$

With simplification (6) the electric potential Φ is depending solely on the electron charge density $\bar{\sigma}_e(\rho)$ and the nuclear charge density $\bar{\sigma}_n(\rho)$. The magnetic vector potential \bar{A}_z is solely depending on the electric current density $\bar{J}_z(\rho)$. The electron charge density and the electric current density are derived from the modulus square of the other electron’s wave functions. This approach has **similarities with the density functional theory (DFT)** used in quantum chemistry modeling, except that the exchange and correlation energies are not accounted for. However, the influence of the magnetic field on the co-linear electron-electron interaction is maintained and the Pauli exclusion principle will be obeyed during orbital occupation.

4.4 Boundary Conditions for Solutions of the Klein-Gordon Equation

Care must be taken according to simplification (20), that **the total energy $\bar{E} + m_e c^2$ of an eigenstate is always positive**, therefore:

$$(33) \quad \bar{E} > -m_e c^2$$

Requirement (33) can be fulfilled by **excluding eigenstates with a negative total energy during orbital occupation**.

According to simplification (22) only bound eigenstates are considered. Therefore the wave function amplitude must disappear at infinite radial distances:

$$(34) \quad \lim_{\rho \rightarrow \infty} \Psi(\rho) = 0$$

As required by simplification (5) the wave function has to meet circular boundary conditions:

$$(35) \quad \Psi(z = 0) = \Psi(z = \bar{L}) \text{ and}$$

$$(36) \quad \frac{\partial \Psi(z = 0)}{\partial z} = \frac{\partial \Psi(z = \bar{L})}{\partial z}$$

By definition of a bound state the total energy of the electron is less than its rest energy:

$$(37) \quad \bar{E} + m_e c^2 < m_e c^2, \text{ therefore } \bar{E} < 0$$

Combining (33) and (37) yields for bound states:

$$(38) \quad -2m_e c^2 < \bar{E} < 0$$

For computing observables the Klein-Gordon electron wave functions have to be **normalized** such that:

$$(39) \quad 1 = \|\Psi\| = \iiint_{\mathbf{R}^3} |\Psi|^2 d\vec{r}$$

4.5 Observables of the Klein-Gordon Electron Wave Function

The potential energy of an electron is solely stemming from the Coulomb field:

$$(40) \quad \bar{E}_{pot}(\vec{r}) = -e\Phi(\vec{r})$$

The local kinetic energy of the electron is, what's left when the potential energy is subtracted from \bar{E} :

$$(41) \quad \bar{E}_{kin}(\vec{r}) = \bar{E} + e\Phi(\vec{r})$$

The **volume charge density distribution of electron number i** in a static electromagnetic potential computes as following:

$$(42) \quad \bar{\sigma}_{e,i} = -e|\Psi_i|^2$$

Summing this up for all N electrons of the CP and using modulus square factorization (69) is resulting in:

$$(43) \quad \bar{\sigma}_e = -e \sum_{i=1}^N |\Psi_i|^2 = -\frac{e}{2\pi\bar{L}} \sum_{i=1}^N |\Psi_{\rho,i}|^2, \text{ where } \Psi_{\rho,i} \text{ is the radial wave function of electron number } i$$

The **current density distribution of electron number i** in a static electromagnetic potential computes as following:

$$(44) \quad \vec{J}_i = \frac{-e}{m_e} \left[-\frac{i\hbar}{2} (\Psi_i^* \nabla \Psi_i - \Psi_i \nabla \Psi_i^*) + e\vec{A} |\Psi_i|^2 \right]$$

Summing (44) up for all N electrons of the CP provides:

$$(45) \quad \vec{J} = \frac{-e}{m_e} \sum_{i=1}^N \left[-\frac{i\hbar}{2} (\Psi_i^* \nabla \Psi_i - \Psi_i \nabla \Psi_i^*) + e\vec{A} |\Psi_i|^2 \right]$$

Using product ansatz (67), modulus square factorization (69) and Ψ_z -solution (73), the **z-component** (in cylindrical coordinates) **of the current density** (45) in a CP computes as:

$$(46) \quad \begin{aligned} \bar{J}_z &= \frac{-e}{m_e} \sum_{i=1}^N \left[-\frac{i\hbar}{2} \left(\Psi_i^* \frac{\partial \Psi_i}{\partial z} - \Psi_i \frac{\partial \Psi_i^*}{\partial z} \right) + e\bar{A}_z |\Psi_i|^2 \right] \\ &= \frac{-e}{2\pi m_e \bar{L}} \sum_{i=1}^N \bar{p}_{z,i} |\Psi_{\rho,i}|^2, \text{ where} \end{aligned}$$

$$(47) \quad \bar{p}_{z,i} = \bar{P}_{z,i} + e\bar{A}_z = \hbar k_i + e\bar{A}_z \text{ is the z-component of the electron's kinetic momentum } \hat{p}$$

When (43), (45) and (46) will be used for determining the electric and magnetic potentials in the Klein-Gordon equation (32), the electron number i is incorrectly exposed also to its own potential. However, this error is quite small, if the CP contains very many electrons.

The total **current** in z-direction carried by all electrons of the CP can be computed by integrating (46) over all radius values and azimuth values:

$$(48) \quad \begin{aligned} I_z &= \int_{\varphi=0}^{2\pi} \int_{\rho=0}^{\infty} \bar{J}_z(\rho) \rho d\rho d\varphi = \frac{-e}{m_e \bar{L}} \int_0^{\infty} \sum_{i=1}^N \bar{p}_{z,i} |\Psi_{\rho,i}(\rho)|^2 \rho d\rho \\ &= \frac{-e}{m_e \bar{L}} \sum_{i=1}^N \int_0^{\infty} [\hbar k_i + e\bar{A}_z(\rho)] |\Psi_{\rho,i}(\rho)|^2 \rho d\rho \end{aligned}$$

The expectation value of the **electron group velocity's z-component** (averaged over all N electrons of the CP) can be computed from the z-component of the total current:

$$(49) \quad \langle v_z \rangle = \frac{I_z \bar{L}}{-Ne}$$

The **expectation value of the electron orbit radius** for eigenstates of equation (32) is:

$$(50) \quad \langle \rho \rangle = \int_0^{\infty} |\Psi|^2 \rho^2 d\rho$$

One could naively assume, that the **total binding energy** \bar{E}_B of a CP is the sum of the energies of all electrons plus the nuclear self-repulsion energy:

$$(51) \quad \bar{E}_B \neq \bar{E}_n + \sum_{i=1}^N \bar{E}_i, \text{ where } \bar{E}_i \text{ is the energy eigenvalue of electron number } i \text{ and } \bar{E}_n \text{ is the nuclear self-repulsion energy defined in (82)}$$

However, this approach would count the electron-electron interaction energies $\langle \bar{E}_{C,e} \rangle_i$ twice, because the electrons are interacting among themselves.

Instead, one needs to subtract half of $\langle \bar{E}_{C,e} \rangle_i$ from the eigenvalues \bar{E}_i before summation:

$$(52) \quad \bar{E}_B = \bar{E}_n + \sum_{i=1}^N \left(\bar{E}_i - \frac{1}{2} \langle \bar{E}_{C,e} \rangle_i \right),$$

where $\langle \bar{E}_{C,e} \rangle_i = \langle -e\Phi_e \rangle_i$ is the expectation value of the electron-electron Coulomb energy for electron number i and Φ_e is the potential of all electrons

4.6 The Electromagnetic Potential and Field of a CP

The electric potential of a CP splits as follows:

$$(53) \quad \Phi = \Phi_n + \Phi_e + \Phi_h, \text{ where } \Phi_n \text{ is the electric potential of the nuclear jellium according to simplification (7) and (8), } \Phi_e \text{ is the electric potential of the electrons and } \Phi_h \text{ is the electric potential of the halo (102).}$$

As a tool for computing the electromagnetic potential the following geometry is analyzed:

A sample charge at distance ρ from the z-axis (origin) and azimuth φ shall act as the point of measurement for vector potential A_z and the electric potentials Φ .

The following figure illustrates this further:

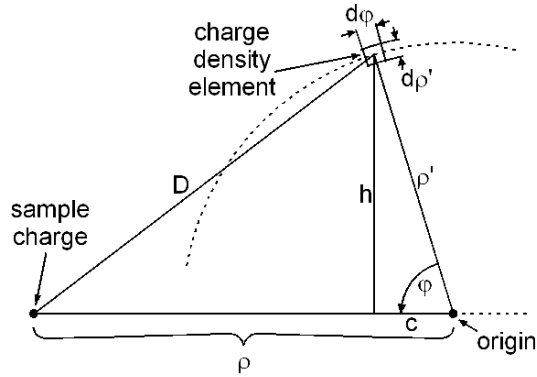


Figure 4 Scheme for computing the electromagnetic potential. This shows a cut perpendicular to the z-axis.

From the geometry of Figure 4 it can be concluded:

$$(54) \quad c = \rho' \cos \varphi$$

$$(55) \quad h = \rho' \sin \varphi$$

$$(56) \quad D = \sqrt{(\rho - c)^2 + h^2} = \sqrt{(\rho - \rho' \cos \varphi)^2 + \rho'^2 \sin^2 \varphi} = \sqrt{\rho'^2 - 2\rho'\rho \cos \varphi + \rho^2}$$

The following figure shall illustrate the geometry in z-direction:

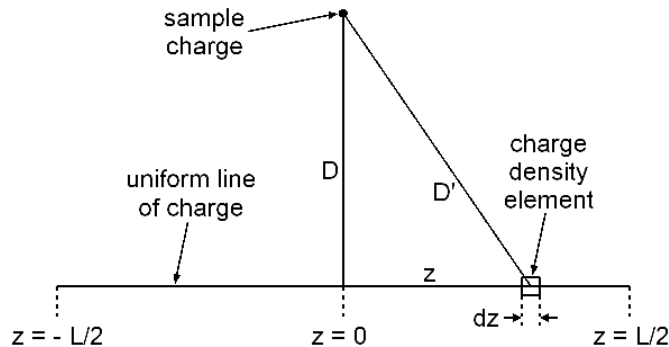


Figure 5 Scheme for computing the electromagnetic potential. This shows a cut in parallel to the z-axis.

From Figure 5 it can be concluded:

$$(57) \quad D' = \sqrt{z^2 + D^2}$$

Figure 5 shows an infinitesimal thin line of charge extending from $z = -\bar{L}/2$ to $z = \bar{L}/2$. This line is in parallel to the z-axis. The volume charge density $\bar{\sigma}(\rho')$ is constant along the line. An infinitesimal charge density element with a volume of $\rho'd\rho'd\varphi dz$ contains a charge of:

$$(58) \quad dQ = \bar{\sigma}(\rho')\rho'd\rho'd\varphi dz, \text{ where}$$

$$(59) \quad \bar{\sigma}(\rho') = \bar{\sigma}_n(\rho') + \bar{\sigma}_e(\rho') + \bar{\sigma}_h(\rho'), \text{ i.e. the sum of the nuclear charge density in the core, the electron charge density and the charge density of the halo (101)}$$

The electric potential at distance D from the infinitesimal line of charge and at axial position $z = 0$ computes as following:

$$(60) \quad d\Phi(D) = \frac{1}{4\pi\epsilon_0} \int_{-\bar{L}/2}^{\bar{L}/2} \frac{dQ}{D'} = \frac{1}{4\pi\epsilon_0} \bar{\sigma}(\rho')\rho'd\rho'd\varphi 2 \int_0^{\bar{L}/2} \frac{1}{\sqrt{z^2 + D^2}} dz$$

$$= \frac{1}{4\pi\epsilon_0} \bar{\sigma}(\rho')\rho'd\rho'd\varphi 2 \ln \frac{\bar{L}/2 + \sqrt{(\bar{L}/2)^2 + D^2}}{D}$$

Inserting (56) into (60) and integrating over ρ' and φ yields the contribution of the entire CP to the **electric potential** (in Lorentz gauge, static case):

$$(61) \quad \Phi(\rho) = \frac{1}{4\pi\epsilon_0} \int_0^\infty \bar{\sigma}(\rho') G(\rho, \rho') \rho' d\rho', \text{ where}$$

$$(62) \quad G(\rho, \rho') = 2 \int_0^{2\pi} \ln \frac{\bar{L}/2 + \sqrt{(\bar{L}/2)^2 + \rho'^2 - 2\rho'\rho \cos \varphi + \rho^2}}{\sqrt{\rho'^2 - 2\rho'\rho \cos \varphi + \rho^2}} d\varphi \text{ is the geometry integral}$$

The geometry integral can be solved by considering Gauss's law and by using the fact that the charge distribution is rotationally symmetric around the z-axis. For the case $\rho \leq \rho'$ there is $G(\rho, \rho') = G(\rho', 0)$. For the case $\rho > \rho'$ there is $G(\rho, \rho') = G(\rho, 0)$:

$$(63) \quad G(\rho, \rho') = \begin{cases} 4\pi \left[\ln \left(\bar{L}/2 + \sqrt{\bar{L}^2/4 + \rho'^2} \right) - \ln \rho' \right] & \text{for } \rho \leq \rho' \\ 4\pi \left[\ln \left(\bar{L}/2 + \sqrt{\bar{L}^2/4 + \rho^2} \right) - \ln \rho \right] & \text{for } \rho > \rho' \end{cases}$$

Replacing $\bar{\sigma}(\rho')/\epsilon_0$ with $\mu_0 \bar{J}_z(\rho')$ in (61) provides the z-component of the CP's **magnetic vector potential** (in Lorentz gauge, static case):

$$(64) \quad \bar{A}_z(\rho) = \frac{\mu_0}{4\pi} \int_0^\infty \bar{J}_z(\rho') G(\rho, \rho') \rho' d\rho' = -\frac{e\mu_0}{8\pi^2 m_e \bar{L}} \int_0^\infty \sum_{i=1}^N [\bar{P}_{z,i} + e\bar{A}_z(\rho')] |\Psi_{\rho',i}|^2 G(\rho, \rho') \rho' d\rho',$$

where I_z is the total current of the electrons in z-direction, μ_0 is the vacuum permeability and $\bar{J}_z(\rho')$ is the z-component of the current density

Note, that \bar{A}_z is depending on itself in equation (64). Therefore, the values of \bar{A}_z and \bar{J}_z need to be determined iteratively until self-consistency.

Based on the circular boundary condition (5) the electric potential (61) and the vector potential (64) are made to be constant in z-direction. This approximation is required for maintaining the full cylindrical symmetry of the model.

The radial and azimuthal (see simplification (15)) components of the vector potential and the current density is zero everywhere. Due to simplification (17) the nuclear jellium is not contributing to the current density.

Note, that $\lim_{\rho \rightarrow \infty} \Phi(\rho) = 0$ and $\lim_{\rho \rightarrow \infty} \bar{A}_z(\rho) = 0$. Equations (61) and (64) therefore can be used for determining the binding energy of electrons to a CP without engaging a non-zero reference potential.

The **radial electric field** of a CP computes as:

$$(65) \quad \varepsilon_\rho = -\frac{\partial}{\partial \rho} \Phi(\rho)$$

The **azimuthal magnetic field** of a CP computes as:

$$(66) \quad B_\varphi = -\frac{\partial}{\partial \rho} \bar{A}_z(\rho)$$

4.7 Product Ansatz

The following **product ansatz** is made to factorize the electron wave function:

$$(67) \quad \Psi(\rho, \varphi, z) = \Psi_\rho(\rho) \Psi_\varphi(\varphi) \Psi_z(z) \text{ or in short: } \Psi = \Psi_\rho \Psi_\varphi \Psi_z$$

The wave function of a single electron is supposed to be normalized and it represents a stationary state. In azimuthal direction and in axial direction the electromagnetic potential is constant. Therefore the modulus square of Ψ_φ and Ψ_z is also constant:

$$(68) \quad |\Psi_\varphi|^2 = \Psi_\varphi^* \Psi_\varphi = \frac{1}{2\pi} \text{ and } |\Psi_z|^2 = \Psi_z^* \Psi_z = \frac{1}{L}$$

Hence the **modulus square** of the entire wave function factorizes as:

$$(69) \quad |\Psi|^2 = \Psi_\rho^*(\rho) \Psi_\rho(\rho) \Psi_\varphi^*(\varphi) \Psi_\varphi(\varphi) \Psi_z^*(z) \Psi_z(z) = \frac{1}{2\pi L} |\Psi_\rho(\rho)|^2$$

The **normalization criteria** (39) could then be carried out as:

$$(70) \quad 1 = \|\Psi\| = \int_0^\infty |\Psi_\rho(\rho)|^2 \rho d\rho$$

4.8 Separation of the Klein-Gordon Equation

With product ansatz (67) the partial derivatives of the wave function are:

$$(71) \quad \frac{\partial \Psi}{\partial \rho} = \Psi_\varphi \Psi_z \frac{d\Psi_\rho}{d\rho} \text{ and } \frac{\partial \Psi}{\partial \varphi} = \Psi_\rho \Psi_z \frac{d\Psi_\varphi}{d\varphi} \text{ and } \frac{\partial \Psi}{\partial z} = \Psi_\rho \Psi_\varphi \frac{d\Psi_z}{dz}$$

Inserting this into the equation (32) and dividing both sides by Ψ yields:

$$(72) \quad -\frac{\hbar^2}{2m_e} \left[\frac{1}{\rho \Psi_\rho} \frac{d}{d\rho} \left(\rho \frac{d\Psi_\rho}{d\rho} \right) + \frac{1}{\rho^2 \Psi_\varphi} \frac{d^2 \Psi_\varphi}{d\varphi^2} + \frac{1}{\Psi_z} \frac{d^2 \Psi_z}{dz^2} + 2 \frac{e \bar{A}_z}{\hbar \Psi_z} i \frac{d\Psi_z}{dz} - \frac{e^2 \bar{A}_z^2}{\hbar^2} \right] - \frac{m_e c^2}{2} \left(\frac{\bar{E} + e\Phi}{m_e c^2} + 1 \right)^2 + \frac{m_e c^2}{2} = 0$$

Remark for the mathematical purity: The division by Ψ is done here out of convenience. It could have been postponed to a later step without affecting the end result, such that wave functions (which can have zeros) never show up in the denominator.

The following **wave function is solving the z-dependent part** of (72):

$$(73) \quad \Psi_z = \sqrt{\frac{1}{L}} e^{ikz}, \text{ where } k \in \mathbf{R}$$

Due to simplification (5) the energy eigenvalues are quantized to a discrete spectrum, because wave number k has to meet the following boundary condition:

$$(74) \quad k = l \frac{2\pi}{L}, \text{ where } l \in \mathbf{Z}$$

Integer l acts as an **axial quantum number** here (This quantum number l should not be confused with the l in Laplace's spherical harmonic function $Y_l^m(\theta, \varphi)$ used for modeling the electrons of atoms).

The following **wave function is solving the φ -dependent part** of (72):

$$(75) \quad \Psi_\varphi = \sqrt{\frac{1}{2\pi}} e^{im\varphi}, \text{ where } m \in \mathbf{Z}$$

Integer m is the **azimuthal quantum number**.

Inserting (47), (73) and (75) into (72) provides the **radial Klein-Gordon equation** of a CP:

$$(76) \quad \left\{ \frac{\hbar^2}{2m_e} \left[-\frac{1}{\rho} \frac{d}{d\rho} \left(\rho \frac{d}{d\rho} \right) + \frac{m^2}{\rho^2} \right] + \frac{p_z^2}{2m_e} - \frac{m_e c^2}{2} \left(\frac{\bar{E} + e\Phi}{m_e c^2} + 1 \right)^2 + \frac{m_e c^2}{2} \right\} \Psi_\rho = 0$$

At the **non-relativistic limit** the term $x = (\bar{E} + e\Phi)/(m_e c^2)$ approaches zero. By using only the first two terms of the Taylor series of $(x+1)^2$ about $x=0$ one can approximate:

$$(77) \quad (x+1)^2 \approx 1 + 2x$$

With this approximation equation (76) becomes the **radial Schrödinger equation** of a CP:

$$(78) \quad \left\{ \frac{\hbar^2}{2m_e} \left[-\frac{1}{\rho} \frac{d}{d\rho} \left(\rho \frac{d}{d\rho} \right) + \frac{m^2}{\rho^2} \right] + \frac{p_z^2}{2m_e} - \bar{E} - e\Phi \right\} \Psi_\rho = 0$$

The radial Schrödinger equation (78) is based on the **non-relativistic Hamiltonian** for an electron in an electromagnetic field with minimal coupling:

$$(79) \quad \hat{H} = \frac{(-i\hbar\nabla + e\vec{A})^2}{2m_e} - e\Phi = \frac{(\vec{P} + e\vec{A})^2}{2m_e} - e\Phi$$

The eigenstates of differential equation (76) or (78) provide the radial wave functions Ψ_ρ . The eigenvalues \bar{E} of bound states are discrete, i.e. they are countable by a principal quantum number n , the azimuthal quantum number m and the axial quantum number l . The **principal quantum number** $n = 1, 2, 3, \dots$ is defined here analogous to the hydrogen atom: n equals one plus the number of node lines of $\Psi_\rho \Psi_\varphi$, therefore $n \geq |m| + 1$ (In a stricter sense, Ψ_φ has no node lines. However, a standing wave of two superposed azimuthal wave functions, differing only in the sign of quantum number m , has m node lines.)

Principal quantum number n has no explicit representation in (76) or (78) or in any of the following formulas. It is useful however, as an ordering scheme for computational results.

One has to keep in mind, that the eigenvalues \bar{E} , the eigenstates Ψ_ρ , Ψ_φ and Ψ_z , as well as the quantum numbers n , m and l are generally distinct for each electron of the CP. In order to ease readability, the electron number as an index has been omitted from these symbols, unless the index is needed in a summation.

4.9 The Jellium Model of the Nuclear Charge Distribution

According to simplification (7) the charge of the nuclei is treated as if it were a uniform "positive jelly" background, rather than point charges with distances in between.

The nuclear charge density distribution $\bar{\sigma}_n(\rho)$ of the jellium has cylindrical symmetry, i.e. it doesn't depend on φ and z . It is a function of the radial distance ρ .

According to equations (53), (59) and (61) the **electric potential of the core nuclear jellium** is:

$$(80) \quad \Phi_n(\rho) = \frac{1}{4\pi\epsilon_0} \int_0^\infty \bar{\sigma}_n(\rho') G(\rho, \rho') \rho' d\rho'$$

An infinitesimal charge element $\bar{\sigma}_n(\rho') \rho' d\rho' d\varphi dz$ brought into potential Φ_n has the potential energy:

$$(81) \quad d\bar{E}_n = \bar{\sigma}_n(\rho) \Phi_n(\rho) \rho d\rho d\varphi dz$$

Integrating (81) over the entire space and dividing the result by two yields the **nuclear self-repulsion energy**:

$$(82) \quad \bar{E}_n = -\bar{E}_{G,n} + \bar{E}_{n,h} + \bar{E}_{h,h} + \frac{1}{2} \int_0^{\bar{L}} \int_0^{2\pi} \int_0^\infty \bar{\sigma}_n(\rho) \Phi_n(\rho) \rho d\rho d\varphi dz$$

$$= -\bar{E}_{G,n} + \bar{E}_{n,h} + \bar{E}_{h,h} + \pi \bar{L} \int_0^\infty \bar{\sigma}_n(\rho) \Phi_n(\rho) \rho d\rho,$$

where $\bar{E}_{G,n}$ is the granularity correction (90), $\bar{E}_n > 0$ and $\bar{E}_{G,n} > 0$, $\bar{E}_{n,h}$ is the halo-core repulsion energy according to equation (104) and $\bar{E}_{h,h}$ is the halo self-repulsion energy according to equation (105)

The division by two in (82) takes care of the fact, that the jellium is interacting with itself and the repulsion energy must not be accounted twice during integration.

According to simplification (8) equation (82) **needs to be corrected by $\bar{E}_{G,n}$ in order to account for the granularity of the nuclear charges**. For this purpose the following approximation is made:

Regarding simplification (13) the nuclei are assumed to have a mean charge of Z_m . The mean charge is determined by the average of the nuclear charges Z_i of the atomic sort weighted by the fraction $0 < F_i < 1$ of the respective atomic sort:

$$(83) \quad Z_m = \sum_0^n F_i Z_i, \text{ where } \sum_0^n F_i = 1 \text{ and } n \text{ is the number of different atomic sorts of the mixture}$$

The **volume occupied by one nucleus** would be:

$$(84) \quad V_{1n} = \frac{eZ_m}{\bar{\sigma}_n} = \frac{4\pi}{3} R_{1n}^3$$

Hence the **radius R_{1n} of a sphere with volume V_{1n}** would be:

$$(85) \quad R_{1n} = \sqrt[3]{\frac{3}{4\pi} \frac{eZ_m}{\bar{\sigma}_n}}, \text{ therefore } \bar{\sigma}_n = \frac{3}{4\pi} \frac{eZ_m}{R_{1n}^3}$$

Assuming that the charge density $\bar{\sigma}_n$ is constant within V_{1n} , the electric potential caused by this sphere is:

$$(86) \quad \Phi_{1n}(r) = \begin{cases} \frac{1}{4\pi\epsilon_0} \frac{eZ_m}{r} & r > R_{1n} \\ \frac{eZ_m}{8\pi\epsilon_0} \left(\frac{3}{R_{1n}} - \frac{r^2}{R_{1n}^3} \right) & r \leq R_{1n} \end{cases}$$

The self-repulsion energy of the jellium within sphere V_{1n} (i.e. for case $r \leq R_{1n}$) would be as following:

$$(87) \quad \begin{aligned} \bar{E}_{1n}(\bar{\sigma}_n) &= \frac{1}{2} \iiint_{V_{1n}} \Phi_{1n}(r) dQ = \frac{1}{2} \int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^{R_{1n}} \bar{\sigma}_n \Phi_{1n}(r) r^2 dr \sin\theta d\theta d\varphi = 2\pi \bar{\sigma}_n \int_0^{R_{1n}} \Phi_{1n}(r) r^2 dr \\ &= \frac{3}{16\pi\epsilon_0} \frac{e^2 Z_m^2}{R_{1n}^3} \int_0^{R_{1n}} \left(\frac{3r^2}{R_{1n}} - \frac{r^4}{R_{1n}^3} \right) dr = \frac{3}{20\pi\epsilon_0} \frac{e^2 Z_m^2}{R_{1n}} = \frac{1}{5\epsilon_0} \sqrt[3]{\left(\frac{3}{4\pi}\right)^2} e^5 Z_m^5 \bar{\sigma}_n \end{aligned}$$

In reality \bar{E}_{1n} , i.e. the self-repulsion energy of a single nucleus, is zero. This is, because a nucleus doesn't repel itself. Therefore the self-repulsion energy (82) needs to be subtracted by (87) for each nucleus in the CP.

An infinitesimal cylindrical zone of a CP with radius $\rho - \frac{1}{2}d\rho \leq \rho' \leq \rho + \frac{1}{2}d\rho$ has the volume:

$$(88) \quad dV = 2\pi\bar{L}\rho d\rho$$

The number of nuclei residing in volume dV is:

$$(89) \quad dN = \frac{dV}{V_{1n}} = 2\pi\bar{L} \frac{\bar{\sigma}_n(\rho)}{eZ_m} \rho d\rho$$

Multiplying (89) with \bar{E}_{1n} and integrating over ρ yields the **granularity correction of the jellium's self-repulsion energy**:

$$(90) \quad \bar{E}_{G,n} = \frac{2\pi\bar{L}}{5\epsilon_0} \left(\frac{3eZ_m}{4\pi} \right)^{2/3} \int_0^\infty \bar{\sigma}_n^{4/3}(\rho) \rho d\rho$$

Similarly, there is a **granularity correction needed for electron-nucleus Coulomb interaction**.

An infinitesimal electron charge element $\bar{\sigma}_e(\rho')\rho'd\rho'd\varphi dz$ brought into the nuclear potential Φ_n of the core has the potential energy:

$$(91) \quad d\bar{E}_e = \bar{\sigma}_e(\rho)\Phi_n(\rho)\rho d\rho d\varphi dz$$

Integrating (91) over the entire space yields the **total electron-nucleus Coulomb energy of the core**:

$$(92) \quad \bar{E}_e = -\bar{E}_{G,e} + \int_0^{\bar{L}} \int_0^{2\pi} \int_0^\infty \bar{\sigma}_e(\rho)\Phi_n(\rho)\rho d\rho d\varphi dz = -\bar{E}_{G,e} + 2\pi\bar{L} \int_0^\infty \bar{\sigma}_e(\rho)\Phi_n(\rho)\rho d\rho, \text{ where } \bar{E}_{G,e} \text{ is the granularity correction (95), } \bar{E}_e < 0 \text{ and } \bar{E}_{G,e} > 0$$

If a sphere with volume V_{1n} is filled with a homogeneous nuclear jellium, the jellium causes a potential as specified in (86). However, if the single nucleus in this sphere is modeled as a point charge (i.e. the reality), then it causes the usual Coulomb potential:

$$(93) \quad \tilde{\Phi}_{1n}(r) = \frac{eZ_m}{4\pi\epsilon_0} \frac{1}{r}, \text{ where } r \leq R_{1n}$$

Assuming that the charge densities $\bar{\sigma}_e$ and $\bar{\sigma}_n$ are constant within V_{1n} , the granularity correction of the electron-nucleus interaction energy per nucleus, \tilde{E}_{1n} , can be computed from the difference between equation (86) and equation (93):

$$\begin{aligned}
(94) \quad \tilde{E}_{1n}(\bar{\sigma}_n, \bar{\sigma}_e) &= \iiint_{V_{1n}} (\Phi_{1n} - \tilde{\Phi}_{1n}) dQ = \int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^{R_{1n}} \bar{\sigma}_e (\Phi_{1n} - \tilde{\Phi}_{1n}) r^2 dr \sin \theta d\theta d\varphi \\
&= 4\pi \bar{\sigma}_e \int_0^{R_{1n}} (\Phi_{1n} - \tilde{\Phi}_{1n}) r^2 dr = \frac{\bar{\sigma}_e}{\epsilon_0} eZ_m \int_0^{R_{1n}} \left(\frac{3r^2}{2R_{1n}} - \frac{r^4}{2R_{1n}^3} - r \right) dr \\
&= -\frac{\bar{\sigma}_e}{10\epsilon_0} eZ_m R_{1n}^2 = -\frac{\bar{\sigma}_e}{10\epsilon_0} (eZ_m)^{5/3} \left(\frac{3}{4\pi\bar{\sigma}_n} \right)^{2/3}
\end{aligned}$$

Multiplying (89) with \tilde{E}_{1n} and integrating over ρ yields the **granularity correction of the electron-nucleus interaction energy in the core**:

$$(95) \quad \bar{E}_{G,e} = -\frac{\pi\bar{L}}{5\epsilon_0} \left(\frac{3eZ_m}{4\pi} \right)^{2/3} \int_0^\infty \bar{\sigma}_e(\rho) \sqrt[3]{\bar{\sigma}_n(\rho)} \rho d\rho$$

In order to make (95) applicable to the Klein-Gordon equation one has to compute it for a single electron by using (69):

$$(96) \quad \bar{E}_{G,e,i} = \frac{e}{10\epsilon_0} \left(\frac{3eZ_m}{4\pi} \right)^{2/3} \int_0^\infty |\Psi_{\rho,i}|^2 \sqrt[3]{\bar{\sigma}_n(\rho)} \rho d\rho$$

Unfortunately, the eigenstate $\Psi_{\rho,i}$ needs to be known before $\bar{E}_{G,e,i}$ can be computed. The latter in turn is required to compute the eigenstate $\Psi_{\rho,i}$ as a solution to the Klein-Gordon equation, providing a vicious cycle. Therefore, $\bar{E}_{G,e,i}$ can only be computed iteratively.

According to simplification (9) the **nucleic charge distribution** is modeled by means of a two-dimensional normal distribution in radial direction:

$$(97) \quad \bar{\sigma}_n(\rho) = \frac{Q}{\bar{L}} \frac{1}{2\pi\bar{s}^2} \exp\left(-\frac{\rho^2}{2\bar{s}^2}\right),$$

where \bar{s} is the standard deviation in meter, Q is the nuclear charge of the entire CP

The distribution function (97) is normalized, such that the integral over all space (in Cartesian coordinates) yields the nuclear charge Q :

$$(98) \quad \iiint \bar{\sigma}_n(\rho) = \int_{x=-\infty}^\infty \int_{y=-\infty}^\infty \int_{z=0}^{\bar{L}} \frac{Q}{\bar{L}} \frac{1}{2\pi\bar{s}^2} \exp\left(-\frac{\rho^2}{2\bar{s}^2}\right) dz dy dx = Q, \text{ where } \rho^2 = x^2 + y^2$$

The standard deviation is to be determined iteratively (by variation), such that the total energy of the CP is minimized. During each computational iteration the electron eigenstates have to be recomputed, as they strongly depend on the nuclear charge distribution.

After studying the computational results of the radial electron density distribution, it became apparent, that equation (97) needs to be adjusted like this:

$$(99) \quad \bar{\sigma}_n(\rho) = \text{Height} \cdot \exp\left(-\frac{\rho^2}{2\bar{s}^2} - \frac{\rho^5}{\text{Slope}^5} - \frac{\rho^{16}}{\text{Cutoff}^{16}}\right)$$

Parameter *Height* has to be computed, such that the distribution function (99) is normalized to the nuclear charge Q . Parameters *Slope* and *Cutoff* should be adjusted, such that the total energy of the CP is minimized.

4.10 The CP Halo (i.e. the Charge Compensation Layer)

The computational results achieved with the cylindrical model are showing, that CP can have an internal chemical potential below -300 keV. The core of a CP can carry excess negative charge (for example, 25% more electrons than nuclear charges).

When a CP is in thermodynamic equilibrium with surrounding matter, the surplus negative charge of the CP has to be compensated by a surrounding layer of cations. This layer contains room charge, which terminates the electrical field around the CP, such that the external chemical potential is flat (i.e. the same in the CP and the surrounding material).

The charge compensation layer of cations will be called the “halo” of the CP, whereas the nuclei and the electrons comprising the CP (without the halo) will be called the “core” of the CP.

The halo can also be modeled as a jellium, like the core.

If the halo is fully compensating the charge of the core, the **linear charge density of the halo** is:

$$(100) \quad \bar{\lambda}_h \equiv \frac{Q_h}{L} = -\bar{\lambda}_e - \bar{\lambda}_n, \text{ where } Q_h \text{ is the total charge of the halo, } \bar{\lambda}_n = Q/L \text{ is the linear charge density of the nuclei in the core and } \bar{\lambda}_e \text{ is the linear charge density of the electrons in the core}$$

For the purpose of computing the electric potential in the core, the easiest way is to assume a homogeneously charged cylindrical shell with a radius ρ_h , which is larger than the radial extent of the electron orbits and the extent of the nuclear charge distribution of the core (simplification (10)). The shell shall be concentric to the z-axis and have an infinitesimal wall thickness of $\delta\rho$. The **charge density of the halo** cations then is:

$$(101) \quad \bar{\sigma}_h = \frac{\bar{\lambda}_h}{2\pi\rho_h\delta\rho}$$

The potential summands sourced by the nuclei and the electrons had been specified in equation (53). The **contribution of the halo to the Coulomb potential** computes as:

$$(102) \quad \Phi_h = \frac{1}{4\pi\epsilon_0} \int_0^\infty \bar{\sigma}_h(\rho') G(\rho, \rho') \rho' d\rho' = \frac{\bar{\lambda}_h}{8\pi^2\epsilon_0} \int_{\rho_h}^{\rho_h+\delta\rho} \frac{1}{\rho'\delta\rho} G(\rho, \rho') \rho' d\rho' = \frac{\bar{\lambda}_h}{8\pi^2\epsilon_0} G(\rho_h, \rho_h)$$

The **contribution of the halo to the Coulomb energy of a single electron** is:

$$(103) \quad \bar{E}_{C,h} \equiv -e\Phi_h = -\frac{e\bar{\lambda}_h}{8\pi^2\epsilon_0} G(\rho_h, \rho_h)$$

The additional charge of the halo increases the self-repulsion energy between the nuclear charges, as mentioned in equation (82). In particular:

The **core-halo repulsion energy** can be computed via:

$$(104) \quad \bar{E}_{n,h} = Q\Phi_h$$

The **halo self-repulsion energy** is:

$$(105) \quad \bar{E}_{h,h} = \frac{1}{2} Q_h \Phi_h$$

The result in the above equation was divided by two, because the halo charges are interacting with themselves.

The halo is the part of the CP, which interchanges matter with the environment. It would therefore be interesting to model the properties (density distribution, electric field, pressure, etc.). Unfortunately, this is not trivial: The internal pressure of the halo is partially caused by electrostatic repulsion of the cations. It is also caused by the degeneracy pressure of the cation's core electron shells. The degeneracy pressure depends on how many electrons remain to be bound to the cations.

Depending on the electrical field strength, part of the cation's electron shells is ripped away. One can therefore expect, that the cations close to the CP core are carrying multiple positive charges, whereas the cations at the outer boundary of the halo are carrying only one positive charge. It was beyond the capacity of the author to model this reliably.

When a CP attaches itself to the surface of a dielectric (or metallic) substrate, some of the cationic charges are replaced by polarization charges induced in the substrate by the electric field of the CP. Per unit of charge it takes less energy to create polarization charge than it takes to create cationic charge. Therefore, a CP can lower its total energy by attaching to a surface and thereby neutralizing some of the halo cations.

4.11 A Plausibility Check on Pressures

Dividing equation (43) by the electron charge provides the **electron density** $d_e = \sum_{i=1}^N |\Psi_i|^2$. By idealizing the electrons as a **free electron gas**, the electron density can be used for computing the **electron degeneracy pressure** (at the non-relativistic limit):

$$(106) \quad P_d = \frac{(3\pi^2)^{2/3} \hbar^2}{5m_e} d_e^{5/3}$$

The degeneracy (outward) pressure is approximately balancing the (inward) pressure generated by the electrostatic and magnetic forces acting on the electrons:

$$(107) \quad P_d \approx P_e + P_m, \text{ where } P_e \text{ is the electrostatic pressure, } P_m \text{ is the magnetic pressure}$$

The above equation is useful as an **plausibility check** for any solutions produced by the Klein-Gordon equation.

The electric field $\varepsilon_\rho(\rho')$ from equation (65) creates a force $F_{\rho,e}(\rho')$ on the infinitesimal charge element $\bar{\sigma}_e(\rho')\rho'd\rho'd\varphi dz$ of the electron gas:

$$(108) \quad F_{\rho,e}(\rho') = \varepsilon_\rho(\rho')\bar{\sigma}_e(\rho')\rho'd\rho'd\varphi dz$$

Integrating (108) provides the **pressure** on the electron gas **created by the electric field**:

$$(109) \quad P_e(\rho) = P_G + \int_\rho^\infty \frac{F_{\rho,e}(\rho')}{\rho'd\varphi dz} = P_G + \int_\rho^\infty \varepsilon_\rho(\rho')\bar{\sigma}_e(\rho')d\rho', \text{ where } P_G \text{ is the granularity correction of the electrostatic pressure described below}$$

The magnetic field $B_\varphi(\rho')$ from equation (66) creates a force $F_{\rho,m}(\rho')$ on the infinitesimal current element $J_z(\rho')\rho'd\rho'd\varphi$ of the electron gas:

$$(110) \quad F_{\rho,m}(\rho') = \bar{L}B_\varphi(\rho')J_z(\rho')\rho'd\rho'd\varphi$$

Integrating (110) provides the **pressure** on the moving electron gas **created by the magnetic field**:

$$(111) \quad P_m(\rho) = \int_\rho^\infty \frac{F_{\rho,m}(\rho')}{\bar{L}\rho'd\varphi} = \int_\rho^\infty B_\varphi(\rho')J_z(\rho')d\rho'$$

The granularity correction P_G stems from electrostatic attraction of the electron gas to the nuclear point charges in the direct vicinity $r \leq R_{1n}$ of the nuclei:

$$(112) \quad P_G(\rho) = \frac{\partial \bar{E}_{en}}{\partial V_{1n}}, \text{ where } \bar{E}_{en} \text{ is the potential energy of the electron gas interacting with the potential of one nucleus } \tilde{\Phi}_{1n} \text{ and with the potential of itself } \Phi'_{1n} \text{ in a sphere with radius } R_{1n}$$

The potential of one nucleus $\tilde{\Phi}_{1n}$ is given by equation (93). The potential of the (homogeneous) electron gas is:

$$(113) \quad \Phi'_{1n}(r) = -\frac{eZ_m \bar{\sigma}_e}{8\pi\epsilon_0 \bar{\sigma}_n} \left(\frac{3}{R_{1n}} - \frac{r^2}{R_{1n}^3} \right), \text{ where } r \leq R_{1n}$$

An infinitesimal electron charge element $dQ = \bar{\sigma}_e(\rho')\rho'd\rho'd\phi dz$ brought into the potential $\tilde{\Phi}_{1n} + \frac{1}{2}\Phi'_{1n}$ has the infinitesimal potential energy:

$$(114) \quad d\bar{E}_e = \bar{\sigma}_e(\rho) \left[\tilde{\Phi}_{1n}(\rho) + \frac{1}{2}\Phi'_{1n}(\rho) \right] \rho d\rho d\phi dz$$

The potential created by the electron gas was divided by two in the above equation, because the electron gas is interacting with itself.

Assuming $\bar{\sigma}_e$ and $\bar{\sigma}_n$ are constant and integrating (114) over the volume V_{1n} provides the potential energy of the electron gas:

$$(115) \quad \begin{aligned} \bar{E}_{en}(\bar{\sigma}_n, \bar{\sigma}_e) &= \iiint_{V_{1n}} \tilde{\Phi}_{1n} dQ = \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^{R_{1n}} \bar{\sigma}_e \left(\tilde{\Phi}_{1n} + \frac{1}{2}\Phi'_{1n} \right) r^2 dr \sin\theta d\theta d\phi \\ &= 4\pi\bar{\sigma}_e \int_0^{R_{1n}} \left(\tilde{\Phi}_{1n} + \frac{1}{2}\Phi'_{1n} \right) r^2 dr = \frac{3}{4\pi\epsilon_0} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \frac{e^2 Z_m^2}{R_{1n}^3} \int_0^{R_{1n}} \left[r - \frac{1}{2} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \left(\frac{3r^2}{R_{1n}} - \frac{r^4}{R_{1n}^3} \right) \right] dr \\ &= \frac{3}{4\pi\epsilon_0} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \frac{e^2 Z_m^2}{R_{1n}} \left(\frac{1}{2} - \frac{2}{5} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \right) = \left(\frac{3}{4\pi} \right)^{2/3} \frac{e^2 Z_m^2}{\epsilon_0} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \left(\frac{1}{2} - \frac{2}{5} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \right) V_{1n}^{-1/3} \end{aligned}$$

Putting (112) and (115) together, using (84) and using the fact, that $\bar{\sigma}_e/\bar{\sigma}_n$ is independent of V_{1n} changes, yields:

$$(116) \quad \begin{aligned} P_G(\rho) &= \frac{\partial \bar{E}_{en}}{\partial V_{1n}} = - \left(\frac{3}{4\pi} \right)^{2/3} \frac{e^2 Z_m^2}{\epsilon_0} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \left(\frac{1}{6} - \frac{2}{15} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \right) V_{1n}^{-4/3} \\ &= - \left(\frac{3eZ_m}{4\pi} \right)^{2/3} \frac{1}{\epsilon_0} \bar{\sigma}_n^{1/3} \bar{\sigma}_e \left(\frac{1}{6} - \frac{2}{15} \frac{\bar{\sigma}_e}{\bar{\sigma}_n} \right) \end{aligned}$$

One has to keep in mind, though, that equation (107) is merely a coarse approximation:

- At low densities, equation (116) overestimates the electrostatic pressure, because the electron gas is strongly polarized around the nuclei. The polarization cloud does not expand so much at volume changes, i.e. the cloud doesn't contribute much to the pressure.
- At high densities, equation (116) underestimates the electrostatic pressure, because the electric field of the nuclei is not yet fully screened by the electron gas at radius R_{1n} . The pressure contribution of the electrostatic attraction between the nuclei and the electron gas at $r > R_{1n}$ has been neglected, though.
- At high densities, equation (106) overestimates the degeneracy pressure P_d , because a fraction of the electrons have a relativistic velocity. At the relativistic limit P_d is proportional to $d_e^{4/3}$, rather than $d_e^{5/3}$ as suggested by (106).

4.12 Transformation to Natural Units

In the following text the **Hartree energy** will be used as a **unit of measure for energy**. It is defined as:

$$(117) \quad \bar{E}_h \equiv \frac{\hbar^2}{m_e a_0^2} = m_e \left(\frac{e^2}{4\pi\epsilon_0 \hbar} \right)^2 = m_e c^2 \alpha^2 \approx 27.211 eV, \text{ where}$$

$$(118) \quad a_0 \equiv \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = \frac{\hbar}{m_e c \alpha} \approx 52.918 pm \text{ is the } \mathbf{Bohr \ radius} \text{ and}$$

$$(119) \quad \alpha \equiv \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c} = \frac{\hbar}{m_e c a_0} \approx 7.2974 \cdot 10^{-3} \text{ is the } \mathbf{fine \ structure \ constant}.$$

The **electron rest energy** in units of \bar{E}_h therefore becomes:

$$(120) \quad \frac{m_e c^2}{\bar{E}_h} = \frac{1}{\alpha^2}$$

The following equation defines a **reference radius**:

$$(121) \quad \rho_0 \equiv \frac{a_0}{\sqrt{\lambda_n}}, \text{ where}$$

$$(122) \quad \lambda_n \equiv \frac{Q}{e} \frac{1}{L} \text{ is the } \mathbf{linear \ nuclear \ charge \ density} \text{ in natural units,}$$

$$(123) \quad L \equiv \frac{\bar{L}}{a_0} \text{ is the } \mathbf{CP \ length} \text{ in units of the Bohr radius.}$$

The definition of the reference radius was crafted, such that the relative radial extent of the electron orbits at the non-relativistic limit becomes independent of the linear nuclear charge density.

The **relative radius** is defined as:

$$(124) \quad r \equiv \frac{\rho}{\rho_0}$$

The **volume charge density** in natural units is defined here as:

$$(125) \quad \sigma = \sigma_n + \sigma_e + \sigma_h \equiv \frac{a_0^3}{e} \bar{\sigma} = \frac{a_0^3}{e} (\bar{\sigma}_n + \bar{\sigma}_e + \bar{\sigma}_h)$$

The **current density** in natural units is defined here as:

$$(126) \quad J_z \equiv \frac{a_0^3}{ec} \bar{J}_z$$

Additionally, the following quantities are defined here:

$$(127) \quad E \equiv \bar{E} / \bar{E}_h, \text{ i.e. the } \mathbf{sum \ of \ the \ potential \ energy \ and \ the \ kinetic \ energy} \text{ of the electron, which is functioning as the } \mathbf{energy \ eigenvalue} \text{ of the Klein-Gordon equation}$$

$$(128) \quad E_C = E_{G,e,i} + E_{C,n} + E_{C,e} + E_{C,h} \equiv \bar{E}_{G,e,i}/\bar{E}_h - e\Phi/\bar{E}_h \\ = \bar{E}_{G,e,i}/\bar{E}_h - e\Phi_n/\bar{E}_h - e\Phi_e/\bar{E}_h - e\Phi_h/\bar{E}_h,$$

i.e. the **potential energy related to the Coulomb potential** of the electrons, the core nuclei and the halo as seen by an electron (negative sample charge)

$$(129) \quad E_{n,h} \equiv \bar{E}_{n,h}/\bar{E}_h = -\lambda_n L E_{C,h}, \text{ i.e. the } \mathbf{core-halo repulsion energy in natural units}$$

$$(130) \quad E_{h,h} \equiv \bar{E}_{h,h}/\bar{E}_h = -\frac{1}{2}\lambda_h L E_{C,h}, \text{ i.e. the } \mathbf{halo self-repulsion energy in natural units}$$

$$(131) \quad P_z \equiv a_0 \bar{P}_z/\hbar = 2\pi l/L \text{ is the } \mathbf{the axial canonical momentum}$$
 of the electron in natural units

$$(132) \quad A_z \equiv -ea_0 \bar{A}_z/\hbar \text{ is the } \mathbf{axial magnetic vector potential}$$
 in natural units. The related terms $-P_z A_z$ and $A_z^2/2$ are the **magnetic electron-electron interaction energy** and the so-called **diamagnetic energy**, respectively

$$(133) \quad p_z \equiv a_0 \bar{p}_z/\hbar = P_z - A_z = 2\pi l/L + ea_0 \bar{A}_z/\hbar \text{ is the } \mathbf{axial kinetic momentum}$$
 of the electron in natural units. The related term $p_z^2/2$ is the **axial kinetic energy** in natural units

$$(134) \quad E_n \equiv \bar{E}_n/\bar{E}_h, \text{ i.e. the } \mathbf{nuclear self repulsion energy}$$

$$(135) \quad E_{G,n} \equiv \bar{E}_{G,n}/\bar{E}_h, \text{ i.e. the } \mathbf{granularity correction}$$
 of the nuclear jellium's self-repulsion

$$(136) \quad E_{G,e,i} \equiv \bar{E}_{G,e,i}/\bar{E}_h, \text{ i.e. the } \mathbf{granularity correction of the electron-nucleus interaction energy}$$

The **radial wave function** in natural units is defined as:

$$(137) \quad R \equiv \rho_0 \Psi_\rho$$

Dividing both sides of (76) by \bar{E}_h , using the product rule of calculus and substituting via (120), (123), (127), (128), (129) and (133) is resulting in:

$$(138) \quad \left\{ -\frac{a_0^2}{2} \frac{d^2}{d\rho^2} - \frac{a_0^2}{2\rho} \frac{d}{d\rho} + \frac{a_0^2}{2} \frac{m^2}{\rho^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left(E - E_C + \frac{1}{\alpha^2} \right)^2 + \frac{1}{2\alpha^2} \right\} \Psi_\rho = 0$$

Substituting (124) and (137) in (138), using the notation R' and R'' for the first and second derivative to r of radial wave function R and multiplying both sides of the equation by ρ_0 yields the **radial Klein-Gordon equation** in natural units:

$$(139) \quad -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[\frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left(E - E_C + \frac{1}{\alpha^2} \right)^2 + \frac{1}{2\alpha^2} \right] R = 0$$

The **Schrödinger equation** (78) in natural units is:

$$(140) \quad -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[\frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} + E_C - E \right] R = 0$$

The **geometry integral** (63) can be expressed in natural units as:

$$(141) \quad G(r, r') = \begin{cases} 4\pi \left[\ln \left(\sqrt{\lambda_n} L / 2 + \sqrt{\lambda_n L^2 / 4 + r'^2} \right) - \ln r' \right] & \text{for } r \leq r' \\ 4\pi \left[\ln \left(\sqrt{\lambda_n} L / 2 + \sqrt{\lambda_n L^2 / 4 + r^2} \right) - \ln r \right] & \text{for } r > r' \end{cases}$$

Multiplying (43) with a_0^3/e and using (69) and (137) yields the **volume charge density in natural units**:

$$(142) \quad \sigma_e(r) = -\frac{\lambda_n}{2\pi L} \sum_{i=1}^N |R_i(r)|^2$$

Multiplying (46) with $a_0^3/(ec)$ and using (74), (119), (123), (129) and (137) yields the **current density in natural units**:

$$(143) \quad J_z(r) = \frac{-\alpha \lambda_n}{2\pi L} \sum_{i=1}^N [P_{z,i} - A_z(r)] |R_i(r)|^2$$

By using (119), (124), (129), (132), (137) in equation (48) the **total current** (averaged over all electrons, in amperes) can be computed from the quantities in natural units as following:

$$(144) \quad I_z = -\frac{e\alpha c}{a_0 L} \sum_{i=1}^N \int_0^\infty [P_{z,i} - A_z(r)] |R_i(r)|^2 r dr$$

Multiplying both sides of (61) by $-e/\bar{E}_h$ and substituting via (124), (125), (128) and (142) provides the **Coulomb energy in natural units**:

$$(145) \quad E_C(r) = -\frac{1}{\lambda_n} \int_0^\infty \sigma(r') G(r, r') r' dr' = E_{C,h} + \frac{1}{2\pi} \int_0^\infty \left\{ \frac{1}{L} \left[\sum_{i=1}^N |R_i(r')|^2 \right] - \frac{\lambda_n}{s^2} \exp\left(-\frac{r'^2}{2s^2}\right) \right\} G(r, r') r' dr'$$

The **contribution of the halo to the Coulomb energy of a single electron in natural units** is:

$$(146) \quad E_{C,h} = -\frac{\lambda_h}{2\pi} G(r_h, r_h), \text{ where } \lambda_h = \frac{a_0 \bar{\lambda}_h}{e} \text{ and } r_h = \frac{\rho_h}{\rho_0}$$

Dividing both sides of equation (96) by \bar{E}_h and substituting via (117), (118), (124), (125) and (137) yields the **granularity correction of the electron-nucleus interaction energy** is:

$$(147) \quad E_{G,e,i} = \frac{(4\pi)^{1/3}}{10} (3Z_m)^{2/3} \int_0^\infty |R_i|^2 \sigma_n^{1/3} r dr$$

Multiplying both sides of (64) by $-ea_0/\hbar$, substituting via (119), (121), (124), (126), (132) and (143) and using $\varepsilon_0 \mu_0 = c^{-2}$ provides the **magnetic vector potential in natural units**:

$$(148) \quad A_z(r) = -\frac{\alpha}{\lambda_n} \int_0^\infty J_z(r') G(r, r') r' dr' = \frac{\alpha^2}{2\pi L} \int_0^\infty \left\{ \sum_{i=1}^N [P_{z,i} - A_z(r')] |R_i(r')|^2 \right\} G(r, r') r' dr'$$

By dividing (50) by ρ_0 and using (121), (124) and (137) the **expectation value of the electron orbit radius in natural units** becomes:

$$(149) \quad \langle r \rangle = \int_0^\infty |R|^2 r^2 dr$$

Dividing (82) by \bar{E}_h and using (90), (117), (121), (125), (128) and (135) and taking care of the fact, that the sample charges are positive, yields the **nuclear self-repulsion energy** in natural units:

$$(150) \quad E_n = -E_{G,n} + E_{n,h} + E_{h,h} - \frac{\pi L}{\lambda_n} \int_0^\infty \sigma_n(r) E_{C,n}(r) r dr, \text{ where } E_n > 0 \text{ and}$$

$$E_{G,n} = \frac{8}{5} \frac{L}{\lambda_n} \int_0^\infty \left(\frac{3\pi^2}{4} \sigma_n^2(r) Z_m \right)^{\frac{2}{3}} r dr \text{ is the } \mathbf{granularity\ error} \text{ in natural units, } E_{G,n} > 0$$

Using (121), (137) and (136) the **normalization criteria** (70) in natural units becomes:

$$(151) \quad 1 = \|R\| = \int_0^\infty |R(r)|^2 r dr$$

The **standard deviation of the nuclear charge distribution** in natural units is:

$$(152) \quad s \equiv \frac{\bar{s}}{\rho_0}$$

Multiplying both sides of (97) with a_0^3/e and substituting via (121), (123), (124) and (152) yields the **nuclear charge distribution** in natural units:

$$(153) \quad \sigma_n(r) = \frac{Q\lambda_n}{2\pi e L s^2} \exp\left(-\frac{r^2}{2s^2}\right)$$

With the modifications made in (99) the nuclear charge distribution reads:

$$(154) \quad \sigma_n(r) = C_H \exp\left(-\frac{r^2}{2s^2} - \frac{r^5}{C_S^5} - \frac{r^{16}}{C_C^{16}}\right), \text{ where } C_H = \frac{a_0^3}{e} \text{Height}, C_S = \frac{\text{Slope}}{\rho_0} \text{ and } C_C = \frac{\text{Cutoff}}{\rho_0}$$

The **charge density of the halo** (derived from equation (101)) in natural units reads:

$$(155) \quad \sigma_h = \frac{\lambda_n \lambda_h}{2\pi r_h \delta r}$$

4.13 Approximate Solution of the Radial Wave Function

The following **ansatz** will be used for approximating the **radial wave function**:

$$(156) \quad R(r) = f(r) \cdot \exp(-\zeta r), \text{ where}$$

$f(r)$ is assumed to be a polynomial and $\zeta \in \mathbf{R}^+$ is a tunable scaling factor.

The radial Klein-Gordon equation (139) has a second solution, which is linear independent of the solution gained by ansatz (156). The second solution would be represented by the following ansatz:

$$(157) \quad R(r) = f(r) \cdot \exp(\zeta r), \text{ where } \zeta \in \mathbf{R}^+$$

However, this second solution and all linear combinations with it were incompatible with boundary condition (34). Therefore, this second solution ansatz will not be used.

The first derivative of the radial wave functions (156) reads:

$$(158) \quad R' = (f' - \zeta f) \cdot \exp(-\zeta r)$$

The second derivative of the radial wave functions is:

$$(159) \quad R'' = (f'' - 2\zeta f' + \zeta^2 f) \cdot \exp(-\zeta r)$$

The value of ζ can be determined by analyzing the asymptotic behavior of the wave function R at $r \rightarrow \infty$:

The electromagnetic potential (and therefore the terms E_C and A_z) become zero, when the radius approaches infinity. Also, the terms proportional to $1/r$ and $1/r^2$ disappear at $r \rightarrow \infty$. The Klein-Gordon equation (139) then simplifies to:

$$(160) \quad -\lambda_n R'' + (P_z^2 - \alpha^2 E^2 - 2E)R = 0$$

Inserting (157) and (159) into (160) leads to:

$$(161) \quad \lambda_n (-f'' + 2\zeta f' - \zeta^2 f) + (P_z^2 - \alpha^2 E^2 - 2E)f = 0$$

Assuming function f can be approximated by a polynomial of finite degree, the function dominates over its derivatives at $r \rightarrow \infty$ and therefore the **exponential scaling factor** is:

$$(162) \quad \zeta = \sqrt{\frac{1}{\lambda_n} (P_z^2 - \alpha^2 E^2 - 2E)}, \text{ where } E < 0$$

The case of $E \geq 0$ does not apply here, because of simplification (22). Only the positive value of the square root is valid here, because of ansatz (156).

Solving (162) for the energy provides:

$$(163) \quad E = \frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2} + \frac{1}{\alpha^2} - \frac{1}{\alpha^2}$$

Only the positive value of the square root is valid here, because of boundary condition (38).

Equation (163) in conjunction with boundary condition (38) has interesting consequences:

$$(164) \quad \frac{|P_z|}{\sqrt{\lambda_n}} < \zeta < \sqrt{\frac{1}{\lambda_n} \left(P_z^2 + \frac{1}{\alpha^2} \right)}, \text{ i.e. the scaling factor } \zeta \text{ is limited by means of the axial canonical momentum at the low end and by means of the rest energy at the high end}$$

At the **non-relativistic limit** the **exponential scaling factor** computes as:

$$(165) \quad \zeta = \sqrt{\frac{1}{\lambda_n} (P_z^2 - 2E)}, \text{ thus } E = \frac{1}{2} (P_z^2 - \lambda_n \zeta^2), \text{ where } \frac{|P_z|}{\sqrt{\lambda_n}} < \zeta < \sqrt{\frac{1}{\lambda_n} \left(P_z^2 + \frac{2}{\alpha^2} \right)}$$

Inserting (156), (158), (159) and (163) into radial Klein-Gordon equation (139) is leading to:

$$(166) \quad -\frac{\lambda_n}{2} f'' + \frac{\lambda_n}{2} \left(2\zeta - \frac{1}{r} \right) f' + \left[\frac{\lambda_n \zeta}{2r} + \frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left(\frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2} + \frac{1}{\alpha^2} - E_C \right)^2 + \frac{1}{2\alpha^2} - \frac{\lambda_n \zeta^2}{2} \right] f = 0$$

For the non-relativistic limit the Schrödinger equation (140) is leading to:

$$(167) \quad -\frac{\lambda_n}{2} f'' + \frac{\lambda_n}{2} \left(2\zeta - \frac{1}{r} \right) f' + \left(\frac{\lambda_n \zeta}{2r} + \frac{\lambda_n m^2}{2r^2} - P_z A_z + \frac{A_z^2}{2} + E_C \right) f = 0$$

Solutions to differential equation (166) or (167) consist of eigenvalues of ζ and eigenstates of polynomial f . These solutions can then be used to compute the eigenvalues of E and eigenstates of R of the radial Klein-Gordon equation (139) or Schrödinger equation (140).

Function $f(r)$ can be approximated by a polynomial of r as following:

$$(168) \quad f(r) \approx \sum_{j=0}^J c_j r^{\beta+j} \quad \text{for } \beta \in \mathbf{N}_0 \text{ and } c_j \in \mathbf{R}$$

The summation runs over a number $J+1$ of terms, depending on the desired accuracy of the approximation (in practice, depending on the configuration of the CP, J needs to be 150 through 700 with 80-bit floating point numbers for “reasonable” accuracy).

The (generally arbitrary) phase of the (generally complex) wave function R is chosen, such that the coefficients c_j become real numbers.

Generally, constants c_j and ζ are depending on quantum numbers n , m and l . For simplicity reasons, this dependency is not reflected in the respective indices of these constants.

The first derivative of (168) reads:

$$(169) \quad f'(r) \approx \sum_{j=0}^J (\beta + j) c_j r^{\beta+j-1}$$

The second derivative of (168) is:

$$(170) \quad f''(r) \approx \sum_{j=0}^J (\beta + j)(\beta + j - 1) c_j r^{\beta+j-2}$$

In equation (166) a number of terms can be approximated by a polynomial of degree P :

$$(171) \quad \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left(\frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2 + \frac{1}{\alpha^2}} - E_C \right)^2 + \frac{1}{2\alpha^2} - \frac{\lambda_n \zeta^2}{2} \approx \sum_{p=0}^P b_p r^p,$$

where $P \leq J - 1$ and $b_p \in \mathbf{R}$

At the non-relativistic limit (171) simplifies to:

$$(172) \quad -P_z A_z + \frac{A_z^2}{2} + E_C \approx \sum_{p=0}^P b_p r^p, \quad \text{where } P \leq J - 1 \text{ and } b_p \in \mathbf{R}$$

Approximations (171) and (172) have a limited convergence radius, no matter how large P is made and how the coefficients are chosen. However, for a given closed interval of radius values the approximations can be made arbitrarily precise by choosing P and the coefficients appropriately. A good approximation accuracy has been achieved with $P = 8$.

A suitable approximation can be found by first determining the range $r_0 \leq r \leq r_p$ of relevant radius values reflecting the radial extent of the electron’s wave function. For example, one can choose r_0 and r_p in such a way, that the electron resides with 99.9% probability between these radii and, at the same time, the range is made as small as possible.

Based on this range, additional nodes r_1 through r_{p-1} need to be determined between r_0 and r_p . The nodes should be chosen, such that the approximation error is minimized (e.g. via Chebyshev nodes). These nodes can then be used e.g. by Newton polynomials for interpolation.

Inserting (168), (169), (170) and (171) into (166) and multiplying both sides with $-2r^2/\lambda_n$ yields:

$$(173) \quad \sum_{j=0}^J (\beta + j)(\beta + j - 1) c_j r^{\beta+j} + (1 - 2\zeta r) \sum_{j=0}^J (\beta + j) c_j r^{\beta+j} - \left[\zeta r + m^2 + \frac{2}{\lambda_n} \sum_{p=0}^P b_p r^{p+2} \right] \sum_{j=0}^J c_j r^{\beta+j} = 0$$

By neglecting the terms with potencies of r higher than $\beta + J$ the result can be written as:

$$(174) \quad \sum_{j=0}^J \left\{ [(\beta + j)^2 - m^2] c_j + \zeta(1 - 2\beta - 2j) c_{j-1} - \frac{2}{\lambda_n} \sum_{p=0}^P b_p c_{j-p-2} \right\} r^{\beta+j} + \varepsilon_{cut} = 0,$$

where $c_i = 0$ for $i < 0$

and ε_{cut} is the cut-off error produced by neglecting potencies of r higher than $\beta + J$

The cut-off error computes as:

$$(175) \quad \varepsilon_{cut} = \zeta(-1 - 2\beta - 2J) c_J r^{\beta+J+1} - \frac{2}{\lambda_n} \sum_{p=0}^P \sum_{j=J+1}^{J+p+2} b_p c_{j-p-2} r^{\beta+j}$$

The left hand side of equation (174) equals zero for all values of r . This can only be true, if the coefficients of $r^{\beta+j}$ fulfill the following equation:

$$(176) \quad [(\beta + j)^2 - m^2] c_j + \zeta(1 - 2\beta - 2j) c_{j-1} - \frac{2}{\lambda_n} \sum_{p=0}^P b_p c_{j-p-2} = 0$$

Analyzing the case $j = 0$ gives:

$$(177) \quad \beta^2 - m^2 = 0, \text{ therefore } \beta = |m|$$

Inserting (177) into (176) yields the **iterative formula for computing the coefficients** from the value of c_0 :

$$(178) \quad c_j = \frac{1}{(2|m|j + j^2)} \left\{ \zeta(2|m| + 2j - 1) c_{j-1} + \frac{2}{\lambda_n} \sum_{p=0}^P b_p c_{j-p-2} \right\}, \text{ where } c_i = 0 \text{ for } i < 0$$

Note, that the coefficients c_j are all proportional to each other. Formula (178) stays the same at the non-relativistic limit.

Equation (175) puts additional requirements on the coefficients c_{J-P} through c_J , which contradict the requirements of equation (178). Therefore, the polynomial approximation of the radial wave function cannot be made precise. Unfortunately, ε_{cut} diverges, if J is made too large. The latter effect is caused by rounding errors in conjunction with finite floating-point number precision.

The approximation error becomes minimal, when the last coefficient c_J is zero, which is the case only for the eigenvalues of ξ . Therefore, this defines a method for determining the eigenvalues.

Alternatively one could determine the eigenvalues by using the original Klein-Gordon equation (139) in conjunction with (163) as a measure of error:

$$(179) \quad \delta(r, \zeta) = -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[\frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left(\frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2 + \frac{1}{\alpha^2}} - E_c \right)^2 + \frac{1}{2\alpha^2} \right] R \\ \approx -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left\{ \frac{\lambda_n m^2}{2r^2} + \sum_{p=0}^P b_p r^p + \frac{\lambda_n \zeta^2}{2} \right\} R$$

At the non-relativistic limit one would use the Schrödinger equation (140) in conjunction with (165) as a measure of error:

$$(180) \quad \delta(r, \zeta) = -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[\frac{\lambda_n m^2}{2r^2} - P_z A_z + \frac{A_z^2}{2} + E_C + \frac{\lambda_n \zeta^2}{2} \right] R$$

$$\approx -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left\{ \frac{\lambda_n m^2}{2r^2} + \sum_{p=0}^P b_p r^p + \frac{\lambda_n \zeta^2}{2} \right\} R$$

The function $\delta(r, \zeta)$ is approaching zero for all values of r only at the energy eigenvalues E or ζ .

The value of coefficient c_0 can be determined from ζ by normalization of the wave function R .

Combining (156), (168) and (177) leads to:

$$(181) \quad R \approx \sum_{j=0}^J c_j r^{|m|+j} \exp(-\zeta r)$$

The normalization condition (151) requires:

$$(182) \quad 1 = \|R\|^2 = \int_0^\infty |R(r)|^2 r dr = \int_0^\infty \left(\sum_{j=0}^J c_j r^{|m|+j} \right)^2 \exp(-2\zeta r) r dr$$

That means, one has to scale all c_j proportionally, such that (182) yields the value 1.

In many cases the cut-off error (175) is so small, that the radial wave function can be computed without having any noticeable distortion from it. Unfortunately, the cut-off error is not always small compared to the amplitude of the “real” wave function. The following plot shall illustrate this:

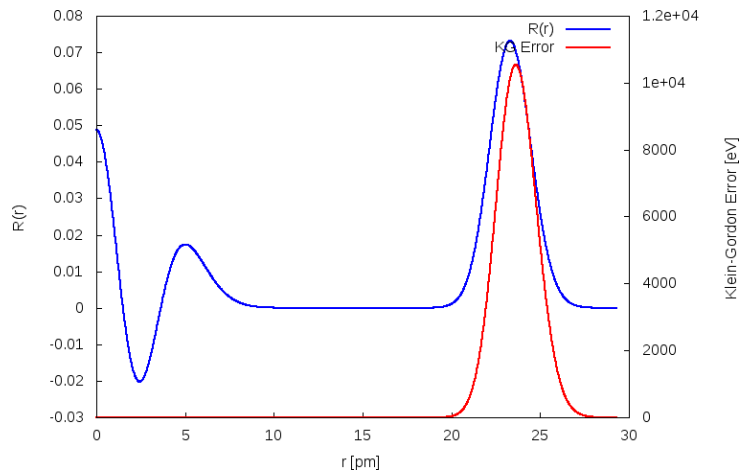


Figure 6 Example of a radial wave function (blue) with a large distortion from the cut-off error centered around $r \approx 23 \text{ pm}$. The distortion is leading to a large Klein-Gordon error (red).

Empirically, the Klein-Gordon error in Figure 6 as computed via (179) is closely correlated to the cut-off error computed via $\varepsilon_{cut} \cdot \exp(-\zeta r)$ for the same wave function:

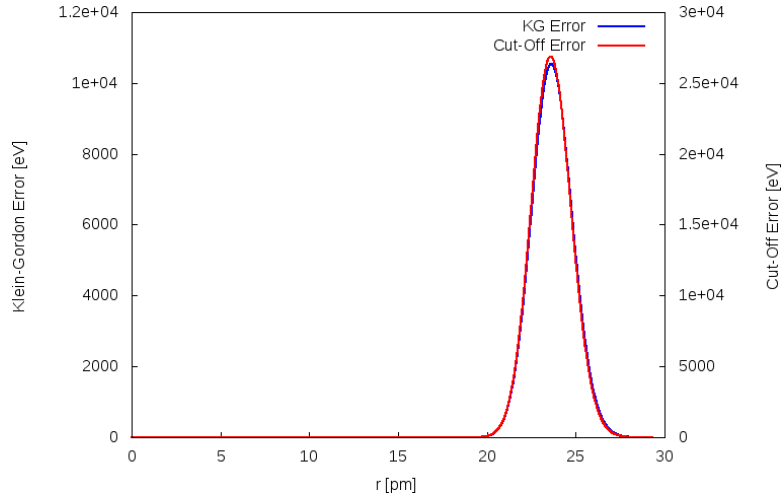


Figure 7 Correlation between the Klein-Gordon error (blue) and the cut-off error (red)

This correlation is a bit surprising, as the Klein-Gordon error is computed from different potencies of the radius than the cut-off error.

Distortions of the wave function and the related Klein-Gordon error are generally becoming smaller, if the upper limits of the summation index in (168) and (169) are adjusted, such that the highest potencies of the radius matches with (170):

$$(183) \quad f = \sum_{j=0}^{J-2} c_j r^{|m|+j} \quad \text{re-definition of (168)}$$

$$(184) \quad f' = \sum_{j=0}^{J-1} (|m| + j) c_j r^{|m|+j-1} \quad \text{re-definition of (169)}$$

$$(185) \quad f'' = \sum_{j=0}^J (|m| + j)(|m| + j - 1) c_j r^{|m|+j-2} \quad \text{copy of (170)}$$

The adjusted upper limits of j are changing the cut-off error (175) to:

$$(186) \quad \varepsilon_{cut} = \left[(|m| + J - 1)^2 c_{J-1} + \zeta (3 - 2|m| - 2J) c_{J-2} \right] r^{|m|+J-1} \\ + (|m| + J - 1) \left[(|m| + J) c_J - 2\zeta c_{J-1} \right] r^{|m|+J} - \frac{2}{\lambda_n} \sum_{p=0}^P \sum_{j=J-1}^{J+p} b_p c_{j-p-2} r^{|m|+j} \quad (\text{Re-definition})$$

4.14 Grouping, Orbital Occupation, Self-Consistent Field Iterations

The electron configuration of a CP consists of many orbitals, which are characterized by the quantum numbers n , m and l . According to the Pauli exclusion principle each orbital can only be occupied by a maximum of two electrons (one with spin up and one with spin down).

There are too many electrons in a CP for computing all occupied orbitals individually. Instead, ranges of orbitals with contiguous values for l are grouped together. Within a group all orbitals have the same quantum numbers n and m . These orbitals of such groups differ in quantum number l . The arithmetic mean of the quantum numbers l represents the group during computation.

The most simple approach is to let each group contain the same number of orbitals. On one hand the groups should be small enough to achieve a fine spacing in the electron energies (for accuracy). On the other hand the groups need to be coarse enough, such that computation time becomes affordable.

Equations (142) and (143) are computed by letting the summation run over the occupied number of groups. Each summand is multiplied by the number of electrons it represents.

For ground state computations the occupation should start with the lowest energy. It should progress to groups with successively higher energy until the targeted number of electrons “found their orbital”.

Equations (142), (143) and (178), as well as the occupation process are depending on each other in a circular manner. Thus they can be computed only **iteratively** until reaching **self-consistency between eigenstates, potential and occupation**.

According to simplification (19) orbitals with eigenvalue energies below $-m_e c^2$ are “forbidden” to occupy. Care must be taken and appropriate numerical damping must be applied, such that fluctuations of orbitals between “forbidden” and “allowed” are not hindering convergence of the SCF algorithm.

Within each of these SCF-iterations (self-consistent field iterations) there is a need for sub-iterations:

According to (143) and (148) the axial current density J_z and the vector potential A_z are mutually depending on each other. Sub-iterations are required for making these quantities consistent with each other, while leaving the eigenstates unchanged.