

### 3 CP Quantum Mechanics

by Lutz Jaitner, May, 2016 through February, 2018

#### 3.1 Basic Assumptions

Modeling of CP is based on the following basic assumptions:

- (1) CP contain ensembles of atomic nuclei densely lined up in a long and very narrow channel.
- (2) The distances between the nuclei are so small, that all electrons bound to these nuclei are delocalized along the channel. In other words: Even in their electronic ground state CP don't consist of individual atoms. CP rather form a quasi-one-dimensional plasma (this could also be seen as a metal).

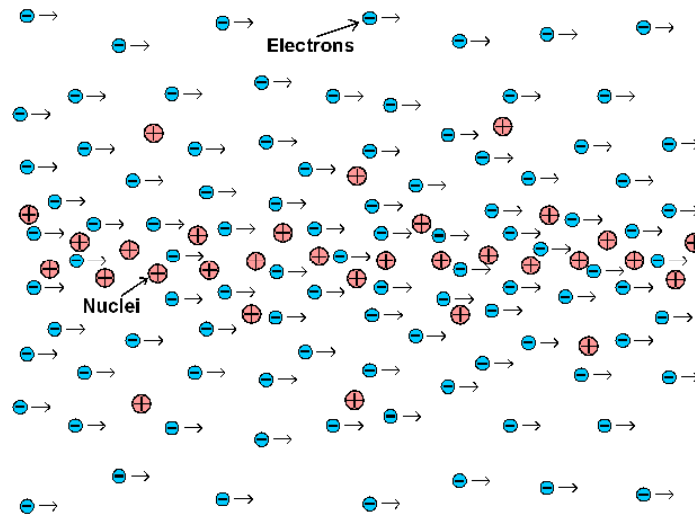


Figure 3 Basic model of a CP. The CP similarly extends to the left and to the right of this picture

#### 3.2 The Cylindrical Model of CP

The shape and quantum mechanical state of CP can be very complicated. In order to obtain a simple quantum mechanical description of CP, the following simplifications are used, which will subsequently be called "the cylindrical model of CP":

- (3) The CP is perfectly straight and cylindrically symmetric, i.e. it is not bent to rings, helices etc. The CP is oriented in parallel to and centered on the z-axis of the modeling cylindrical coordinate system.
- (4) The CP has the length  $\bar{L}$  and contains a total nuclear charge  $Q$  in its core zone (explained below).
- (5) The electron wave functions of the CP are confined in the interval  $0 \leq z < \bar{L}$ . At  $z = \bar{L}$  these wave function are continuously extended to their value and gradient at  $z = 0$ , as if the CP were rings. However, this is meant to describe only the *circular boundary condition* of the wave functions at  $z = \bar{L}$ , not the shape of the CP.
- (6) No external field is applied to the CP.
- (7) A *jellium* model is used for the spatial distribution of the nuclear charge. This means, the positive charges of the nuclei are modeled as a uniform "positive jelly" background, rather than point charges with distances in between. The nucleic charge density is assumed to be constant in axial and azimuthal direction, but it depends on the radial distance.
- (8) The jellium is modeled differently in the core zone and in the halo zone. The narrow core zone is within the reach of the electron wave functions. The diffuse halo zone is outside the reach of the electron wave functions.
- (9) The nucleic charge distribution of the core zone is modeled by means of a two-dimensional normal distribution in radial direction. The fraction of nuclei residing in the core zone, as well as the standard deviation, are to be determined by variation, such that the total energy of the CP is minimized.

- (10) It is assumed, that the halo zone consists of ions each having one positive elementary charge. The ion charge distribution of the halo zone is modeled, such that the positions of the ions are in equilibrium with the electric potential of the CP.
- (11) The CP is assumed to reside in a vacuum. Interaction of the CP with surrounding matter is thus neglected.
- (12) Only stationary states are modeled, as the goal is to describe the ground state of CP. Consequently, the model assumes there is no electron scattering, i.e. there is no momentum transfer between electrons and the nuclei.
- (13) For computing the repulsion energy among the nuclei, short-range corrections to the jellium model have to be made, which account for the granularity of the nuclear charges. In case the CP contains a mixture of different sorts of atomic nuclei, only the mean nuclear charge is taken into account for the corrections in the core zone, rather than the individual nuclear charges.
- (14) The time-independent Klein-Gordon equation is used for modeling the electron wave functions, thereby neglecting the magnetic moments of electron spins. The Klein-Gordon equation is taking care of the large relativistic effects occurring in CP, e.g. the mass defect stemming from the very high binding energy of the electrons. (Clearly, the Dirac equation would be more adequate for modeling CP. However, the involved complexities of such approach are avoided here.). For comparing the formulas and simulation results with the ones obtained from a non-relativistic Hamiltonian, also the Schrödinger equation is used.
- (15) The magnetic field of the azimuthal electron orbits is neglected.
- (16) Magnetic field from nuclear spins is neglected.
- (17) The electron wave functions are modeled in an inertial frame of reference, where no magnetic field is created by any collinear movements of the nuclei. This simplification amounts to an approximation in cases where the nucleic velocities are position dependent.
- (18) The multi-electron system is approximated by computing a collection of one-electron orbitals, whereby each electron orbital is subjected to the mean electric potential and magnetic vector potential created by the total charge density and total current density of all other occupied orbitals and the nuclei (independent particle model). The Pauli exclusion principle is used for determining orbital occupations of the ground state. Exchange and correlation energies are neglected.
- (19) Quantum field theory is not engaged. Particle count is conserved.
- (20) Eigenstates are excluded from occupation, where the corresponding total energy eigenvalue (including the electron's rest energy) of the electron is negative. This shall ensure that the mass defect per electron doesn't exceed the electrons rest energy.
- (21) Only bound eigenstates of the electrons are considered, i.e. the total energy of an eigenstate has to be less than the electron rest energy (i.e. the sum of the potential energy and the kinetic energy has to be negative).

### 3.3 The Klein-Gordon Equation of a CP

Initial calculations of a CP with the Schrödinger equation have shown, that the resulting binding energy of the electrons would potentially far exceed the rest energy of the electrons. Also, the axial electron velocities would potentially far exceed the speed of light. These results were absurd in light of special relativity, because the mass defect per bound electron should never exceed twice its rest mass and velocities should always stay lower than the speed of light. This provided ample reasons for engaging a relativistic Hamiltonian and a Lorentz-covariant quantum mechanical equation to model CP.

Generally, the Dirac equation is regarded as the correct Lorentz-covariant equation for modeling fermions, especially when the effects resulting from the particle's spin is of concern. Unfortunately, the Dirac equation involves 4-component wave functions and the solution of four coupled differential equations, resulting in sizeable mathematical and computational efforts.

Assuming that the electron spins have only minor effects on the binding energy, charge density, current density and other observables, the Klein-Gordon equation provides a Lorentz-covariant alternative to the Dirac equation for modeling the electrons of CP. At the non-relativistic limit the Klein-Gordon equation is equivalent to the Schrödinger equation, while both equations share the deficiency of not modeling the spin.

In relativistic electrodynamics with so-called minimal coupling the Hamiltonian (total energy) of a particle with charge  $q$  moving in the presence of a static (external) electromagnetic potential is:

$$(22) \quad \hat{H} = c\sqrt{(\vec{P} - q\vec{A})^2 + m_e^2 c^2} + q\Phi, \text{ where } c \text{ is the speed of light, } \Phi \text{ is the electric potential, } \vec{A} \text{ is the magnetic vector potential, } \vec{P} = \gamma m_e \vec{v} + q\vec{A} \text{ is the electron's canonical momentum, } m_e \text{ is the electron rest mass and } \gamma \text{ is the Lorentz factor}$$

By defining  $\bar{E} \equiv \hat{H} - m_e c^2$  as being the total energy minus the rest energy and by using  $q = -e$  as the charge of an electron, (22) is leading to the following equation for an electron in a static electromagnetic potential:

$$(23) \quad \bar{E} = c\sqrt{(\vec{P} + e\vec{A})^2 + m_e^2 c^2} - m_e c^2 - e\Phi, \text{ where } e \text{ is the elementary charge}$$

Therefore:

$$(24) \quad (\bar{E} + e\Phi + m_e c^2)^2 = (c\vec{P} + e\vec{A})^2 + m_e^2 c^4$$

All formulas are written in SI units, unless otherwise noted. Throughout this document, energy symbols with a bar on top (e.g.  $\bar{E}$ ) denote, that the energy is measured in Joule. Energy symbols without a bar on top denote, that the energy is measured in units of the Hartree energy (97), i.e. the energy is a dimensionless quantity in the respective formula. Likewise, other symbols with a bar (e.g.  $\bar{\sigma}$ ,  $\bar{J}_z$ ,  $\bar{p}_z$ ,  $\bar{P}_z$ ,  $\bar{A}_z$ ) are in SI units, while its counterparts without the bar are in natural units (i.e. dimensionless).

By quantizing the canonical momentum via the del operator  $\vec{P} \equiv -i\hbar\nabla$  and applying both sides to an electron wave function  $\Psi$ , equation (24) transforms to the **stationary Klein-Gordon equation of an electron in a static electromagnetic potential**:

$$(25) \quad (\bar{E} + e\Phi + m_e c^2)^2 \Psi = \left[ (-i\hbar c \nabla + e\vec{A})^2 + m_e^2 c^4 \right] \Psi, \text{ where}$$

$\hbar$  is the reduced Planck constant and  $i = \sqrt{-1}$

Due to simplification (19),  $\Psi$  is called here a “wave function”, rather than a “quantum field”.

The term  $\bar{E} + m_e c^2$  represents the total energy of the electron. Usually the Klein-Gordon equation is written, such that the total energy is sought as the eigenvalue of this differential equation. However, this document deviates from the customary approach. Instead, the quantity  $\bar{E}$  is sought here as the eigenvalue (both approaches are equivalent in their results).

Some authors prefer the term “relativistic Schrödinger equation” for (25), insisting that the Klein-Gordon equation is different. Here, these terms are used interchangeably.

In quantum mechanics a multi-electron system is correctly described by a single wave function  $\Psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N)$  depending on the positions of the  $N$  electrons. The multi-electron wave function is usually formed by a Slater determinant (or a linear combination of several Slater determinants) to ensure anti-symmetry and the Pauli exclusion principle.

However, the number of electrons in a CP can exceed  $10^{10}$ , which renders a Slater determinant entirely impractical to compute, because a program cannot handle equations with e.g.  $10^{10}$  positions and compute determinants of this size.

According to simplification (18) a rigorously simpler approach is used here for modeling CP, requiring only moderate compute power:

So, instead of using a multi-electron Klein-Gordon equation describing the pair-wise interaction between  $N$  electrons, the cylindrical model uses  **$N$  single-electron Klein-Gordon equations with  $N$  wave functions  $\Psi(\vec{r})$** , each describing a single electron in the **mean potential** of all other electrons and the nuclei.

Of course, this is merely an independent particle approximation. For example, the approach doesn't account for the exchange energy and the correlation energy usually deemed important in quantum chemistry.

At first glance this looks still challenging to compute, because there are  $N$  Klein-Gordon equations to be solved. Fortunately, large numbers of these equations can be computed in groups, because they produce nearly the same charge density distributions and current density distributions.

Expanding the right side of (25) and using  $\nabla \cdot \vec{A} = 0$  (Lorentz gauge in the static case) yields:

$$(26) \quad (\bar{E} + e\Phi + m_e c^2)^2 \Psi = \left( -\hbar^2 c^2 \nabla^2 - 2i\hbar e c^2 \vec{A} \cdot \nabla + e^2 c^2 \vec{A} \cdot \vec{A} + m_e^2 c^4 \right) \Psi$$

A kinetic momentum operator is defined here as:

$$(27) \quad \hat{p} = -i\hbar \nabla + e\vec{A}$$

The expectation value of  $\hat{p}$  equals  $\langle \gamma \rangle m_e \langle \vec{v} \rangle$ , hence the name “kinetic momentum” ( $\langle \vec{v} \rangle$  is the expectation value of the electron’s group velocity and  $\langle \gamma \rangle$  is the expectation value of the local Lorentz factor).

Using (27) in (26) yields:

$$(28) \quad (\bar{E} + e\Phi + m_e c^2)^2 \Psi = \left( c^2 \hat{p}^2 + m_e^2 c^4 \right) \Psi, \text{ where } c^2 \hat{p}^2 = -\hbar^2 c^2 \nabla^2 - 2i\hbar e c^2 \vec{A} \cdot \nabla + e^2 c^2 \vec{A} \cdot \vec{A}$$

According to simplification (14) and (15) the magnetic field of the electron spins and of the azimuthal movement of the electrons is neglected. Thus the only source of the magnetic field is the current carried by the electrons moving in z-direction. Therefore, the vector potential is everywhere oriented in z-direction:

$$(29) \quad \vec{A} = \bar{A}_z \vec{e}_z$$

The Laplace operator expands in cylindrical coordinates as following:

$$(30) \quad \nabla^2 = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2}, \text{ where } \rho \text{ is the radial distance from the z-axis, } \varphi \text{ is the azimuth and } z \text{ is the coordinate of the z-axis}$$

Inserting (29) and (30) into equation (26) and dividing both sides by  $2m_e c^2$  is resulting in **the stationary Klein-Gordon equation of an electron in the mean potential of a CP’s all other electrons and the nuclei**:

$$(31) \quad \left\{ \frac{-\hbar^2}{2m_e} \left[ \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2}{\partial \varphi^2} + \frac{\partial^2}{\partial z^2} + 2 \frac{e\bar{A}_z}{\hbar} i \frac{\partial}{\partial z} - \frac{e^2 \bar{A}_z^2}{\hbar^2} \right] - \frac{m_e c^2}{2} \left( \frac{\bar{E} + e\Phi}{m_e c^2} + 1 \right)^2 + \frac{m_e c^2}{2} \right\} \Psi = 0$$

With simplification (6) the electric potential  $\Phi$  is depending solely on the electron charge density  $\bar{\sigma}_e(\rho)$  and the nuclear charge density  $\bar{\sigma}_n(\rho)$ . The magnetic vector potential  $\bar{A}_z$  is solely depending on the electric current density  $\bar{J}_z(\rho)$ . The electron charge density and the electric current density are derived from the modulus square of the other electron’s wave functions. This approach has **similarities with the density functional theory (DFT)** used in quantum chemistry modeling, except that the exchange and correlation energies are not accounted for. However, the influence of the magnetic field on the co-linear electron-electron interaction is maintained and the Pauli exclusion principle will be obeyed during orbital occupation.

### 3.4 Boundary Conditions for Solutions of the Klein-Gordon Equation

Care must be taken according to simplification (20), that **the total energy  $\bar{E} + m_e c^2$  of an eigenstate is always positive**, therefore:

$$(32) \quad \bar{E} > -m_e c^2$$

Requirement (32) can be fulfilled by **excluding eigenstates with a negative total energy during orbital occupation**.

According to simplification (21) only bound eigenstates are considered. Therefore the wave function amplitude must disappear at infinite radial distances:

$$(33) \quad \lim_{\rho \rightarrow \infty} \Psi(\rho) = 0$$

As required by simplification (5) the wave function has to meet circular boundary conditions:

$$(34) \quad \Psi(z = 0) = \Psi(z = \bar{L}) \text{ and}$$

$$(35) \quad \frac{\partial \Psi(z = 0)}{\partial z} = \frac{\partial \Psi(z = \bar{L})}{\partial z}$$

By definition of a bound state the total energy of the electron is less than its rest energy:

$$(36) \quad \bar{E} + m_e c^2 < m_e c^2, \text{ therefore } \bar{E} < 0$$

Combining (32) and (36) yields for bound states:

$$(37) \quad -2m_e c^2 < \bar{E} < 0$$

For computing observables the Klein-Gordon electron wave functions have to be **normalized** such that:

$$(38) \quad 1 = \|\Psi\| = \iiint_{\mathbf{R}^3} |\Psi|^2 d\vec{r}$$

### 3.5 Observables of the Klein-Gordon Electron Wave Function

The potential energy of an electron is solely stemming from the Coulomb field:

$$(39) \quad \bar{E}_{pot}(\vec{r}) = -e\Phi(\vec{r})$$

The local kinetic energy of the electron is, what's left when the potential energy is subtracted from  $\bar{E}$  :

$$(40) \quad \bar{E}_{kin}(\vec{r}) = \bar{E} + e\Phi(\vec{r})$$

The **volume charge density distribution of electron number**  $i$  in a static electromagnetic potential computes as following:

$$(41) \quad \bar{\sigma}_{e,i} = -e|\Psi_i|^2$$

Summing this up for all electrons of the CP is resulting in:

$$(42) \quad \bar{\sigma}_e = -e \sum_{i=1}^N |\Psi_i|^2$$

The **current density distribution of electron number**  $i$  in a static electromagnetic potential computes as following:

$$(43) \quad \vec{J}_i = \frac{-e}{m_e} \left[ -\frac{i\hbar}{2} (\Psi_i^* \nabla \Psi_i - \Psi_i \nabla \Psi_i^*) + e\vec{A} |\Psi|^2 \right]$$

Summing (43) up for all  $N$  electrons of the CP provides:

$$(44) \quad \vec{J} = \frac{-e}{m_e} \sum_{i=1}^N \left[ -\frac{i\hbar}{2} (\Psi_i^* \nabla \Psi_i - \Psi_i \nabla \Psi_i^*) + e\vec{A} |\Psi|^2 \right]$$

Using product ansatz (62), modulus square factorization (64) and  $\Psi_z$ -solution (68), the **z-component** (in cylindrical coordinates) **of the current density** (44) in a CP computes as:

$$(45) \quad \bar{J}_z = \frac{-e}{m_e} \sum_{i=1}^N \left[ -\frac{i\hbar}{2} \left( \Psi_i^* \frac{\partial \Psi_i}{\partial z} - \Psi_i \frac{\partial \Psi_i^*}{\partial z} \right) + e\bar{A}_z |\Psi_i|^2 \right]$$

$$= \frac{-e}{2\pi m_e \bar{L}} \sum_{i=1}^N \bar{p}_{z,i} |\Psi_{\rho,i}|^2, \text{ where}$$

$$(46) \quad \bar{p}_{z,i} = \bar{P}_{z,i} + e\bar{A}_z = \hbar k_i + e\bar{A}_z \text{ is the } \mathbf{z}\text{-component of the electron's kinetic momentum } \hat{p}$$

When (42), (44) and (45) will be used for determining the electric and magnetic potentials in the Klein-Gordon equation (31), the electron number  $i$  is incorrectly exposed also to its own potential. However, this error is quite small, if the CP contains very many electrons.

The total **current** in z-direction carried by all electrons of the CP can be computed by integrating (45) over all radius values and azimuth values:

$$(47) \quad I_z = \int_{\varphi=0}^{2\pi} \int_{\rho=0}^{\infty} \bar{J}_z(\rho) \rho d\rho d\varphi = \frac{-e}{m_e \bar{L}} \int_0^{\infty} \sum_{i=1}^N \bar{p}_{z,i} |\Psi_{\rho,i}(\rho)|^2 \rho d\rho$$

$$= \frac{-e}{m_e \bar{L}} \sum_{i=1}^N \int_0^{\infty} [\hbar k_i + e\bar{A}_z(\rho)] |\Psi_{\rho,i}(\rho)|^2 \rho d\rho$$

The expectation value of the **electron group velocity's z-component** (averaged over all  $N$  electrons of the CP) can be computed from the z-component of the total current:

$$(48) \quad \langle v_z \rangle = \frac{I_z \bar{L}}{-Ne}$$

The z-component of the **local group velocity** of an electron can be computed from (46):

$$(49) \quad v_z = \frac{\bar{p}_z}{\gamma m_e} = \frac{\hbar k + e\bar{A}_z}{\gamma m_e}$$

The **expectation value of the electron orbit radius** for eigenstates of equation (31) is:

$$(50) \quad \langle \rho \rangle = \int_0^{\infty} \gamma(\rho) |\Psi|^2 \rho^2 d\rho$$

### 3.6 The Electromagnetic Potential of a CP

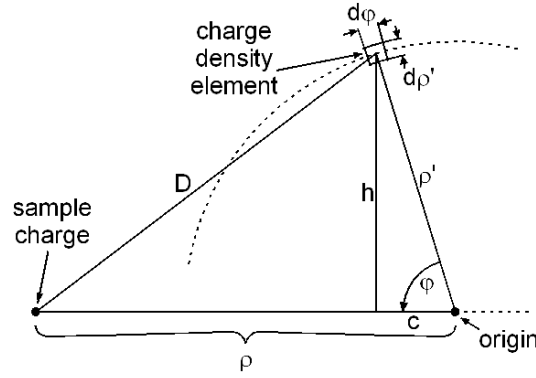
The electric potential of a CP splits as follows:

$$(51) \quad \Phi = \Phi_n + \Phi_e, \text{ where } \Phi_n \text{ is the electric potential of the nuclear jellium according to simplification (7) and } \Phi_e \text{ is the electric potential of the electrons}$$

As a tool for computing the electromagnetic potential the following geometry is analyzed:

A sample charge at distance  $\rho$  from the z-axis (origin) and azimuth  $\varphi$  shall act as the point of measurement for vector potential  $A_z$  and the electric potentials  $\Phi_n$  and  $\Phi_e$ .

The following figure illustrates this further:



**Figure 4** Scheme for computing the electromagnetic potential. This shows a cut perpendicular to the z-axis.

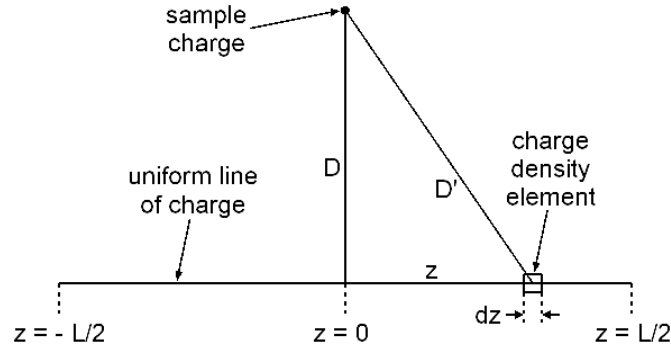
From the geometry of Figure 4 it can be concluded:

$$(52) \quad c = \rho' \cos \varphi$$

$$(53) \quad h = \rho' \sin \varphi$$

$$(54) \quad D = \sqrt{(\rho - c)^2 + h^2} = \sqrt{(\rho - \rho' \cos \varphi)^2 + \rho'^2 \sin^2 \varphi} = \sqrt{\rho'^2 - 2\rho'\rho \cos \varphi + \rho^2}$$

The following figure shall illustrate the geometry in z-direction:



**Figure 5** Scheme for computing the electromagnetic potential. This shows a cut in parallel to the z-axis.

From Figure 5 it can be concluded:

$$(55) \quad D' = \sqrt{z^2 + D^2}$$

Figure 5 shows an infinitesimal thin line of charge extending from  $z = -\bar{L}/2$  to  $z = \bar{L}/2$ . This line is in parallel to the z-axis. The volume charge density  $\bar{\sigma}(\rho')$  is constant along the line. An infinitesimal charge density element with a volume of  $\rho' d\rho' d\varphi dz$  contains a charge of:

$$(56) \quad dQ = \bar{\sigma}(\rho') \rho' d\rho' d\varphi dz, \text{ where}$$

$$(57) \quad \bar{\sigma}(\rho') = [\bar{\sigma}_n(\rho') + \bar{\sigma}_e(\rho')], \text{ i.e. the sum of the nuclear charge density and the electron charge density}$$

The electric potential at distance D from the infinitesimal line of charge and at axial position  $z = 0$  computes as following:

$$(58) \quad d\Phi(D) = \frac{1}{4\pi\epsilon_0} \int_{-\bar{L}/2}^{\bar{L}/2} \frac{dQ}{D'} = \frac{1}{4\pi\epsilon_0} \bar{\sigma}(\rho') \rho' d\rho' d\varphi 2 \int_0^{\bar{L}/2} \frac{1}{\sqrt{z^2 + D^2}} dz$$

$$= \frac{1}{4\pi\epsilon_0} \bar{\sigma}(\rho') \rho' d\rho' d\varphi 2 \ln \frac{\bar{L}/2 + \sqrt{(\bar{L}/2)^2 + D^2}}{D}$$

Inserting (54) into (58) and integrating over  $\rho'$  and  $\varphi$  yields the contribution of the entire CP to the **electric potential** (in Lorentz gauge, static case):

$$(59) \quad \Phi(\rho) = \frac{1}{4\pi\epsilon_0} \int_0^\infty \bar{\sigma}(\rho') G(\rho, \rho') \rho' d\rho', \text{ where}$$

$$(60) \quad G(\rho, \rho') = 2 \int_0^{2\pi} \ln \frac{\bar{L}/2 + \sqrt{(\bar{L}/2)^2 + \rho'^2 - 2\rho'\rho \cos \varphi + \rho^2}}{\sqrt{\rho'^2 - 2\rho'\rho \cos \varphi + \rho^2}} d\varphi \text{ is the geometry integral}$$

Replacing  $\bar{\sigma}(\rho')/\epsilon_0$  with  $\mu_0 \bar{J}_z(\rho')$  in (59) provides the z-component of the CP's **magnetic vector potential** (in Lorentz gauge, static case):

$$(61) \quad \bar{A}_z(\rho) = \frac{\mu_0}{4\pi} \int_0^\infty \bar{J}_z(\rho') G(\rho, \rho') \rho' d\rho' = -\frac{e\mu_0}{8\pi^2 m_e \bar{L}} \int_0^\infty \sum_{i=1}^N [\bar{P}_{z,i} + e\bar{A}_z(\rho')] |\Psi_{\rho,i}|^2 G(\rho, \rho') \rho' d\rho',$$

where  $I_z$  is the total current of the electrons in z-direction,  $\mu_0$  is the vacuum permeability and  $\bar{J}_z(\rho')$  is the z-component of the current density

Note, that  $\bar{A}_z$  is depending on itself in equation (61). Therefore, the values of  $\bar{A}_z$  and  $\bar{J}_z$  need to be determined iteratively until self-consistency.

Based on the circular boundary condition (5) the electric potential (59) and the vector potential (61) are made to be constant in z-direction. This approximation is required for maintaining the full cylindrical symmetry of the model.

The radial and azimuthal (see simplification (15)) components of the vector potential and the current density is zero everywhere. Due to simplification (17) the nuclear jellium is not contributing to the current density.

Note, that  $\lim_{\rho \rightarrow \infty} \Phi(\rho) = 0$  and  $\lim_{\rho \rightarrow \infty} \bar{A}_z(\rho) = 0$ . Equations (59) and (61) therefore can be used for determining the binding energy of electrons to a CP without engaging a non-zero reference potential.

### 3.7 Product Ansatz

The following **product ansatz** is made to factorize the electron wave function:

$$(62) \quad \Psi(\rho, \varphi, z) = \Psi_\rho(\rho) \Psi_\varphi(\varphi) \Psi_z(z) \text{ or in short: } \Psi = \Psi_\rho \Psi_\varphi \Psi_z$$

The wave function of a single electron is supposed to be normalized and it represents a stationary state. In azimuthal direction and in axial direction the electromagnetic potential is constant. Therefore the modulus square of  $\Psi_\varphi$  and  $\Psi_z$  is also constant:

$$(63) \quad |\Psi_\varphi|^2 = \Psi_\varphi^* \Psi_\varphi = \frac{1}{2\pi} \text{ and } |\Psi_z|^2 = \Psi_z^* \Psi_z = \frac{1}{L}$$

Hence the **modulus square** of the entire wave function factorizes as:

$$(64) \quad |\Psi|^2 = \Psi_\rho^*(\rho) \Psi_\rho(\rho) \Psi_\varphi^*(\varphi) \Psi_\varphi(\varphi) \Psi_z^*(z) \Psi_z(z) = \frac{1}{2\pi L} |\Psi_\rho(\rho)|^2$$

The **normalization criteria** (38) could then be carried out as:

$$(65) \quad 1 = \|\Psi\| = \int_0^\infty |\Psi_\rho(\rho)|^2 \rho d\rho$$



### 3.8 Separation of the Klein-Gordon Equation

With product ansatz (62) the partial derivatives of the wave function are:

$$(66) \quad \frac{\partial \Psi}{\partial \rho} = \Psi_\varphi \Psi_z \frac{d\Psi_\rho}{d\rho} \quad \text{and} \quad \frac{\partial \Psi}{\partial \varphi} = \Psi_\rho \Psi_z \frac{d\Psi_\varphi}{d\varphi} \quad \text{and} \quad \frac{\partial \Psi}{\partial z} = \Psi_\rho \Psi_\varphi \frac{d\Psi_z}{dz}$$

Inserting this into the equation (31) and dividing both sides by  $\Psi$  yields:

$$(67) \quad -\frac{\hbar^2}{2m_e} \left[ \frac{1}{\rho \Psi_\rho} \frac{d}{d\rho} \left( \rho \frac{d\Psi_\rho}{d\rho} \right) + \frac{1}{\rho^2 \Psi_\varphi} \frac{d^2 \Psi_\varphi}{d\varphi^2} + \frac{1}{\Psi_z} \frac{d^2 \Psi_z}{dz^2} + 2 \frac{e\bar{A}_z}{\hbar \Psi_z} i \frac{d\Psi_z}{dz} - \frac{e^2 \bar{A}_z^2}{\hbar^2} \right] - \frac{m_e c^2}{2} \left( \frac{\bar{E} + e\Phi}{m_e c^2} + 1 \right)^2 + \frac{m_e c^2}{2} = 0$$

Remark for the mathematical purity: The division by  $\Psi$  is done here out of convenience. It could have been postponed to a later step without affecting the end result, such that wave functions (which can have zeros) never show up in the denominator.

The following **wave function is solving the z-dependent part** of (67):

$$(68) \quad \Psi_z = \sqrt{\frac{1}{L}} e^{ikz}, \quad \text{where } k \in \mathbf{R}$$

Due to simplification (5) the energy eigenvalues are quantized to a discrete spectrum, because wave number  $k$  has to meet the following boundary condition:

$$(69) \quad k = l \frac{2\pi}{L}, \quad \text{where } l \in \mathbf{Z}$$

Integer  $l$  acts as an **axial quantum number** here (This quantum number  $l$  should not be confused with the  $l$  in Laplace's spherical harmonic function  $Y_l^m(\theta, \varphi)$  used for modeling the electrons of atoms).

The following **wave function is solving the  $\varphi$ -dependent part** of (67):

$$(70) \quad \Psi_\varphi = \sqrt{\frac{1}{2\pi}} e^{im\varphi}, \quad \text{where } m \in \mathbf{Z}$$

Integer  $m$  is the **azimuthal quantum number**.

Inserting (46), (68) and (70) into (67) provides the **radial Klein-Gordon equation** of a CP:

$$(71) \quad \left\{ \frac{\hbar^2}{2m_e} \left[ -\frac{1}{\rho} \frac{d}{d\rho} \left( \rho \frac{d}{d\rho} \right) + \frac{m^2}{\rho^2} \right] + \frac{p_z^2}{2m_e} - \frac{m_e c^2}{2} \left( \frac{\bar{E} + e\Phi}{m_e c^2} + 1 \right)^2 + \frac{m_e c^2}{2} \right\} \Psi_\rho = 0$$

At the **non-relativistic limit** the term  $x = (\bar{E} + e\Phi)/(m_e c^2)$  approaches zero. By using only the first two terms of the Taylor series of  $(x+1)^2$  about  $x=0$  one can approximate:

$$(72) \quad (x+1)^2 \approx 1 + 2x$$

With this approximation equation (71) becomes the **radial Schrödinger equation** of a CP:

$$(73) \quad \left\{ \frac{\hbar^2}{2m_e} \left[ -\frac{1}{\rho} \frac{d}{d\rho} \left( \rho \frac{d}{d\rho} \right) + \frac{m^2}{\rho^2} \right] + \frac{p_z^2}{2m_e} - \bar{E} - e\Phi \right\} \Psi_\rho = 0$$

The radial Schrödinger equation (73) is based on the **non-relativistic Hamiltonian** for an electron in an electromagnetic field with minimal coupling:

$$(74) \quad \hat{H} = \frac{(-i\hbar\nabla + e\vec{A})^2}{2m_e} - e\Phi = \frac{(\vec{P} + e\vec{A})^2}{2m_e} - e\Phi$$

The eigenstates of differential equation (71) or (73) provide the radial wave functions  $\Psi_\rho$ . The eigenvalues  $\bar{E}$  of bound states are discrete, i.e. they are countable by a principal quantum number  $n$ , the azimuthal quantum number  $m$  and the axial quantum number  $l$ . The **principal quantum number**  $n = 1, 2, 3, \dots$  is defined here analogous to the hydrogen atom:  $n$  equals one plus the number of node lines of  $\Psi_\rho \Psi_\varphi$ , therefore  $n \geq |m| + 1$  (In a stricter sense,  $\Psi_\varphi$  has no node lines. However, a standing wave of two superposed azimuthal wave functions, differing only in the sign of quantum number  $m$ , has  $m$  node lines.)

Principal quantum number  $n$  has no explicit representation in (71) or (73) or in any of the following formulas. It is useful however, as an ordering scheme for computational results.

One has to keep in mind, that the eigenvalues  $\bar{E}$ , the eigenstates  $\Psi_\rho$ ,  $\Psi_\varphi$  and  $\Psi_z$ , as well as the quantum numbers  $n$ ,  $m$  and  $l$  are generally distinct for each electron of the CP. In order to ease readability, the electron number as an index has been omitted from these symbols, unless the index is needed in a summation.

### 3.9 The Jellium Model of the Nuclear Charge Distribution

According to simplification (7) the charge of the nuclei is treated as if it were a uniform "positive jelly" background, rather than point charges with distances in between.

The nuclear charge density distribution  $\sigma_n(\rho)$  of the jellium has cylindrical symmetry, i.e. it doesn't depend on  $\varphi$  and  $z$ . It is a function of the radial distance  $\rho$ .

According to equations (51), (57) and (59) the **electric potential of the nuclear jellium** is:

$$(75) \quad \Phi_n(\rho) = \frac{1}{4\pi\epsilon_0} \int_0^\infty \bar{\sigma}_n(\rho') G(\rho, \rho') \rho' d\rho'$$

An infinitesimal charge density element  $\bar{\sigma}_n(\rho') \rho' d\rho' d\varphi dz$  brought into potential  $\Phi_n$  has the potential energy:

$$(76) \quad d\bar{E}_n = \bar{\sigma}_n(\rho) \Phi_n(\rho) \rho d\rho d\varphi dz$$

Integrating (76) over the entire space and dividing the result by two yields the **nuclear self-repulsion energy**:

$$(77) \quad \bar{E}_n = -\bar{E}_G + \frac{1}{2} \int_0^{\bar{L}} \int_0^{2\pi} \int_0^\infty \bar{\sigma}_n(\rho) \Phi_n(\rho) \rho d\rho d\varphi dz = -\bar{E}_G + \pi\bar{L} \int_0^\infty \bar{\sigma}_n(\rho) \Phi_n(\rho) \rho d\rho, \text{ where}$$

$\bar{E}_G$  is the granularity correction (85),  $\bar{E}_n > 0$  and  $\bar{E}_G > 0$

The division by two in (77) takes care of the fact, that the jellium is interacting with itself and the repulsion energy must not be accounted twice during integration.

Equation (77) needs to be corrected by  $\bar{E}_G$  in order to account for the granularity of the nuclear charges. For this purpose the following approximation is made:

Regarding simplification (13) the nuclei are assumed to have a mean charge of  $Z_m$ . The mean charge is determined by the average of the nuclear charges  $Z_i$  of the atomic sort weighted by the fraction  $0 < F_i < 1$  of the respective atomic sort:

$$(78) \quad Z_m = \sum_0^n F_i Z_i, \text{ where } \sum_0^n F_i = 1 \text{ and } n \text{ is the number of different atomic sorts of the mixture}$$

The volume occupied by one nucleus would be:

$$(79) \quad V_1 = \frac{eZ_m}{\bar{\sigma}_n} = \frac{4\pi}{3} R^3$$

Hence the radius  $R$  of a sphere with volume  $V_1$  would be:

$$(80) \quad R = \sqrt[3]{\frac{3}{4\pi} \frac{eZ_m}{\bar{\sigma}_n}}$$

Assuming that the charge density  $\bar{\sigma}_n$  is constant within  $V_1$ , the electric potential of the sphere is:

$$(81) \quad \Phi_1(r) = \begin{cases} \frac{\bar{\sigma}_n R^3}{3\epsilon_0 r} = \frac{1}{4\pi\epsilon_0} \frac{eZ_m}{r} & r > R \\ \frac{\bar{\sigma}_n}{6\epsilon_0} (3R^2 - r^2) = \frac{\bar{\sigma}_n}{6\epsilon_0} \left[ 3 \left( \frac{3eZ_m}{4\pi\bar{\sigma}_n} \right)^{\frac{2}{3}} - r^2 \right] & r \leq R \end{cases}$$

The self-repulsion energy of the jellium within sphere  $V_1$  (i.e. for case  $r \leq R$ ) would be as following:

$$(82) \quad \begin{aligned} \bar{E}_1(\bar{\sigma}_n) &= \frac{1}{2} \iiint_V \Phi_1(r) dQ = \frac{1}{2} \int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^R \bar{\sigma}_n \Phi_1(r) r^2 dr \sin\theta d\theta d\varphi = 2\pi\bar{\sigma}_n \int_0^R \Phi_1(r) r^2 dr \\ &= 2\pi\bar{\sigma}_n \int_0^R \frac{\bar{\sigma}_n}{6\epsilon_0} (3R^2 - r^2) r^2 dr = \frac{\pi\bar{\sigma}_n^2}{3\epsilon_0} \left( 3R^2 \int_0^R r^2 dr - \int_0^R r^4 dr \right) \\ &= \frac{1}{5\epsilon_0} \frac{4\pi}{3} \bar{\sigma}_n^2 R^5 = \frac{1}{5\epsilon_0} \sqrt[3]{\left( \frac{3}{4\pi} \right)^2} \bar{\sigma}_n e^5 Z_m^5 \end{aligned}$$

In reality  $\bar{E}_1$ , i.e. the self-repulsion energy of a single nucleus, is zero. This is, because a nucleus doesn't repel itself. Therefore the self-repulsion energy (77) needs to be subtracted by (82) for each nucleus in the CP.

An infinitesimal cylindrical zone of a CP with radius  $\rho_1 - \frac{1}{2}d\rho \leq \rho \leq \rho_1 + \frac{1}{2}d\rho$  has the volume:

$$(83) \quad dV = 2\pi\bar{L}\rho d\rho$$

The number of nuclei residing in volume  $dV$  is:

$$(84) \quad dN = \frac{dV}{V_1} = 2\pi\bar{L} \frac{\bar{\sigma}_n(\rho)}{Z_m} \rho d\rho$$

Multiplying (84) with  $\bar{E}_1$  and integrating over  $\rho$  yields the **granularity correction** of the core jellium's self-repulsion energy:

$$(85) \quad \bar{E}_G = \frac{2\pi\bar{L}}{5\epsilon_0} \int_0^\infty \left( \frac{3\bar{\sigma}_n^2(\rho)eZ_m}{4\pi} \right)^{\frac{2}{3}} \rho d\rho$$

According to simplification (10) the jellium's charge distribution of the **halo zone**  $\bar{\sigma}_{n,h}$  shall be modeled, such that the positions of the cations are in equilibrium with the electric potential of the CP. This is the equilibrium of two pressures resulting from Coulomb forces: The repulsive pressure between neighboring ions and the attractive pressure from the CP's electric potential (59).

Each ion has the following radius:

$$(86) \quad R = \sqrt[3]{\frac{3}{4\pi} \frac{e}{\bar{\sigma}_{n,h}}}$$

A sphere with this radius has a surface area of:

$$(87) \quad A = 4\pi R^2$$

The repulsive Coulomb force between two neighbor ions is:

$$(88) \quad |\vec{F}| = \frac{1}{4\pi\epsilon_0} \frac{e^2}{R^2}$$

As an approximation it is assumed that the Coulomb force component perpendicular to the surface area (87) is constant over this area.

Therefore, the Coulomb force is resulting in a pressure of:

$$(89) \quad P_1 = \frac{|\vec{F}|}{A} = \frac{1}{16\pi^2\epsilon_0} \frac{e^2}{R^4} = \frac{1}{\epsilon_0} \sqrt[3]{\frac{e^2}{\pi^2} \frac{\bar{\sigma}_{n,h}^4}{6^4}}$$

An infinitesimal cylindrical zone (83) contains the charge:

$$(90) \quad dQ = 2\pi\bar{L} \bar{\sigma}_{n,h} \rho d\rho$$

This charge is attracted by the CP's electric potential (59) with the following force:

$$(91) \quad |\vec{F}| = dQ|\vec{E}| = dQ|-\nabla\Phi(\rho)| = \frac{\bar{L} \bar{\sigma}_{n,h}(\rho) \rho d\rho}{2\epsilon_0} \frac{d}{d\rho} \int_0^\infty \bar{\sigma}(\rho') G(\rho, \rho') \rho' d\rho'$$

Force (91) is perpendicular to the surface  $2\pi\bar{L}$  of the cylindrical zone (83), thus creating an incremental pressure of:

$$(92) \quad dP_2(\rho) = \frac{|\vec{F}|}{2\pi\bar{L}} = \frac{1}{4\pi\epsilon_0} \bar{\sigma}_{n,h}(\rho) \rho d\rho \frac{d}{d\rho} \int_0^\infty \bar{\sigma}(\rho') G(\rho, \rho') \rho' d\rho'$$

Integrating (92) over the radius yields the pressure at radial distance  $\rho''$ :

$$(93) \quad P_2(\rho'') = \frac{1}{4\pi\epsilon_0} \int_{\rho''}^\infty \bar{\sigma}_{n,h}(\rho) \left[ \frac{d}{d\rho} \int_0^\infty \bar{\sigma}(\rho') G(\rho, \rho') \rho' d\rho' \right] \rho d\rho$$

Pressures  $P_1$  and  $P_2$  have to be identical in order to maintain an equilibrium of the repulsion between neighboring ions and the global attraction of the jellium by the electric potential of the CP, as is required by simplification (10).

Therefore:

$$(94) \quad \sqrt[3]{\frac{\pi e^2}{162} \bar{\sigma}_{n,h}^4(\rho'')} = \frac{1}{2} \int_{\rho''}^\infty \bar{\sigma}_{n,h}(\rho) \left[ \frac{d}{d\rho} \int_0^\infty \bar{\sigma}(\rho') G(\rho, \rho') \rho' d\rho' \right] \rho d\rho$$

The ion charge density distribution of the halo zone has to be computed iteratively (by variation), such that equation (94) becomes approximately true for all values  $\rho''$  of the halo zone. During these computational iterations the electron wave functions should be kept constant, because there is presumably very little dependency between the electron charge distribution and the nuclear charge distribution of the halo.

According to simplification (9) the **nucleic charge distribution of the core zone** is modeled by means of a two-dimensional normal distribution in radial direction:

$$(95) \quad \bar{\sigma}_{n,c}(\rho) = \frac{Q_c}{L} \frac{1}{2\pi\bar{s}^2} \exp\left(-\frac{\rho^2}{2\bar{s}^2}\right),$$

where  $\bar{s}$  is the standard deviation in meter,  $Q_c$  is the nuclear charge in the core zone

The distribution function (95) is normalized, such that the integral over all space (in Cartesian coordinates) yields the nuclear charge  $Q_c$  of the core:

$$(96) \quad \iiint \bar{\sigma}_{n,c}(\rho) = \int_{x=-\infty}^{\infty} \int_{y=-\infty}^{\infty} \int_{z=0}^{\bar{L}} \frac{Q_c}{L} \frac{1}{2\pi\bar{s}^2} \exp\left(-\frac{\rho^2}{2\bar{s}^2}\right) dz dy dx = Q_c, \text{ where } \rho^2 = x^2 + y^2$$

The fraction of nuclei residing in the core zone, as well as the standard deviation, are to be determined iteratively (by variation), such that the total energy of the CP is minimized. During each computational iteration the electron eigenstates have to be recomputed, as they strongly depend on the nuclear charge distribution of the core.

### 3.10 Transformation to Natural Units

In the following text the **Hartree energy** will be used as a **unit of measure for energy**. It is defined as:

$$(97) \quad \bar{E}_h \equiv \frac{\hbar^2}{m_e a_0^2} = m_e \left( \frac{e^2}{4\pi\epsilon_0 \hbar} \right)^2 = m_e c^2 \alpha^2 \approx 27.211 eV, \text{ where}$$

$$(98) \quad a_0 \equiv \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = \frac{\hbar}{m_e c \alpha} \approx 52.918 pm \text{ is the } \mathbf{Bohr \ radius} \text{ and}$$

$$(99) \quad \alpha \equiv \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c} = \frac{\hbar}{m_e c a_0} \approx 7.2974 \cdot 10^{-3} \text{ is the } \mathbf{fine \ structure \ constant}.$$

The **electron rest energy** in units of  $\bar{E}_h$  therefore becomes:

$$(100) \quad \frac{m_e c^2}{\bar{E}_h} = \frac{1}{\alpha^2}$$

The following equation defines a **reference radius**:

$$(101) \quad \rho_0 \equiv \frac{a_0}{\sqrt{\lambda_n}}, \text{ where}$$

$$(102) \quad \lambda_n \equiv \frac{Q}{e L} \text{ is the } \mathbf{linear \ nuclear \ charge \ density} \text{ in natural units,}$$

$$(103) \quad Q = Q_c + Q_h \text{ is the total nuclear charge, } Q_c \text{ is the nuclear charge in the core zone, } Q_h \text{ is the nuclear charge in the halo zone, and}$$

$$(104) \quad L \equiv \frac{\bar{L}}{a_0} \text{ is the } \mathbf{CP \ length} \text{ in units of the Bohr radius.}$$

The definition of the reference radius was crafted, such that the relative radial extent of the electron orbits at the non-relativistic limit becomes independent of the linear nuclear charge density.

The **relative radius** is defined as:

$$(105) \quad r \equiv \frac{\rho}{\rho_0}$$

The **volume charge density** in natural units is defined here as:

$$(106) \quad \sigma = \sigma_n + \sigma_e \equiv \frac{a_0^3}{e} \bar{\sigma} = \frac{a_0^3}{e} (\bar{\sigma}_n + \bar{\sigma}_e)$$

The **current density** in natural units is defined here as:

$$(107) \quad J_z \equiv \frac{a_0^3}{ec} \bar{J}_z$$

Additionally, the following quantities are defined here:

$$(108) \quad E \equiv \bar{E}/\bar{E}_h, \text{ i.e. the } \mathbf{sum\ of\ the\ potential\ energy\ and\ the\ kinetic\ energy} \text{ of the electron, which is functioning as the } \mathbf{energy\ eigenvalue} \text{ of the Klein-Gordon equation}$$

$$(109) \quad E_C = E_{C,n} + E_{C,e} \equiv -e\Phi/\bar{E}_h = -e\Phi_n/\bar{E}_h - e\Phi_e/\bar{E}_h, \text{ i.e. the } \mathbf{potential\ energy\ related\ to\ the\ Coulomb\ potential} \text{ of the electrons and the nuclei as seen by an electron (negative sample charge)}$$

$$(110) \quad P_z \equiv a_0 \bar{P}_z/\hbar = 2\pi l/L \text{ is the } \mathbf{the\ axial\ canonical\ momentum} \text{ of the electron in natural units}$$

$$(111) \quad A_z \equiv -ea_0 \bar{A}_z/\hbar \text{ is the } \mathbf{axial\ magnetic\ vector\ potential} \text{ in natural units. The related terms } -P_z A_z \text{ and } A_z^2/2 \text{ are the } \mathbf{magnetic\ electron-electron\ interaction\ energy} \text{ and the so-called } \mathbf{diamagnetic\ energy}, \text{ respectively}$$

$$(112) \quad p_z \equiv a_0 \bar{p}_z/\hbar = P_z - A_z = 2\pi l/L + ea_0 \bar{A}_z/\hbar \text{ is the } \mathbf{axial\ kinetic\ momentum} \text{ of the electron in natural units. The related term } p_z^2/2 \text{ is the } \mathbf{axial\ kinetic\ energy} \text{ in natural units}$$

$$(113) \quad E_n \equiv \bar{E}_n/\bar{E}_h, \text{ i.e. the } \mathbf{nuclear\ self\ repulsion\ energy}$$

$$(114) \quad E_G \equiv \bar{E}_G/\bar{E}_h, \text{ i.e. the } \mathbf{granularity\ error} \text{ of the nuclear jellium}$$

The **radial wave function** in natural units is defined as:

$$(115) \quad R \equiv \rho_0 \Psi_\rho$$

Dividing both sides of (71) by  $\bar{E}_h$ , using the product rule of calculus and substituting via (100), (104), (108), (109), (110) and (112) is resulting in:

$$(116) \quad \left\{ -\frac{a_0^2}{2} \frac{d^2}{d\rho^2} - \frac{a_0^2}{2\rho} \frac{d}{d\rho} + \frac{a_0^2}{2} \frac{m^2}{\rho^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left( E - E_C + \frac{1}{\alpha^2} \right)^2 + \frac{1}{2\alpha^2} \right\} \Psi_\rho = 0$$

Substituting (105) and (115) in (116), using the notation  $R'$  and  $R''$  for the first and second derivative to  $r$  of radial wave function  $R$  and multiplying both sides of the equation by  $\rho_0$  yields the **radial Klein-Gordon equation** in natural units:

$$(117) \quad -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[ \frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left( E - E_C + \frac{1}{\alpha^2} \right)^2 + \frac{1}{2\alpha^2} \right] R = 0$$

The **Schrödinger equation** (73) in natural units is:

$$(118) \quad -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[ \frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} + E_C - E \right] R = 0$$

The **geometry integral** (60) can be expressed in natural units as:

$$(119) \quad G(r, r') = 2 \int_0^{2\pi} \ln \frac{\sqrt{\lambda_n L/2 + \sqrt{\lambda_n L^2/4 + r'^2 - 2r'r \cos \varphi + r^2}}}{\sqrt{r'^2 - 2r'r \cos \varphi + r^2}} d\varphi$$

Multiplying (42) with  $a_0^3/e$  and using (64) and (115) yields the **volume charge density in natural units**:

$$(120) \quad \sigma_e(r) = -\frac{\lambda_n}{2\pi L} \sum_{i=1}^N |R_i(r)|^2$$

Multiplying (45) with  $a_0^3/(ec)$  and using (69), (99), (104), (110) and (115) yields the **current density in natural units**:

$$(121) \quad J_z(r) = -\frac{\alpha \lambda_n}{2\pi L} \sum_{i=1}^N [P_{z,i} - A_z(r)] |R_i(r)|^2$$

By using (99), (105), (110), (111), (115) in equation (47) the **total current** (averaged over all electrons, in amperes) can be computed from the quantities in natural units as following:

$$(122) \quad I_z = -\frac{e\alpha c}{a_0 L} \sum_{i=1}^N \int_0^\infty [P_{z,i} - A_z(r)] |R_i(r)|^2 r dr$$

Multiplying both sides of (59) by  $-e/\bar{E}_h$  and substituting via (105), (106), (109) and (120) provides the **Coulomb energy** in natural units:

$$(123) \quad E_C(r) = -\frac{1}{\lambda_n} \int_0^\infty \sigma(r') G(r, r') r' dr' = \frac{1}{2\pi} \int_0^\infty \left\{ \frac{1}{L} \left[ \sum_{i=1}^N |R_i(r')|^2 \right] - \frac{\lambda_n}{s^2} \exp\left(-\frac{r'^2}{2s^2}\right) \right\} G(r, r') r' dr'$$

Multiplying both sides of (61) by  $-ea_0/\hbar$ , substituting via (99), (101), (105), (107), (111) and (121) and using  $\epsilon_0 \mu_0 = c^{-2}$  provides the **magnetic vector potential** in natural units:

$$(124) \quad A_z(r) = -\frac{\alpha}{\lambda_n} \int_0^\infty J_z(r') G(r, r') r' dr' = \frac{\alpha^2}{2\pi L} \int_0^\infty \left\{ \sum_{i=1}^N [P_{z,i} - A_z(r')] |R_i(r')|^2 \right\} G(r, r') r' dr'$$

By dividing (50) by  $\rho_0$  and using (101), (105) and (115) the **expectation value of the electron orbit radius** in natural units becomes:

$$(125) \quad \langle r \rangle = \int_0^\infty |R|^2 r^2 dr$$

Dividing (77) by  $\bar{E}_h$  and using (85), (97), (101), (106), (109) and (114) and taking care of the fact, that the sample charges are positive, yields the **nuclear self-repulsion energy** in natural units:

$$(126) \quad E_n = -E_G - \frac{\pi L}{\lambda_n} \int_0^\infty \sigma_n(r) E_{C,n}(r) r dr, \text{ where } E_n > 0 \text{ and}$$

$$E_G = \frac{8}{5} \frac{L}{\lambda_n} \int_0^\infty \left( \frac{3\pi^2}{4} \sigma_n^2(r) Z_m \right)^{\frac{2}{3}} r dr \text{ is the } \mathbf{granularity error} \text{ in natural units, } E_G > 0$$

Using (101), (115) and 0 the **normalization criteria** (65) in natural units becomes:

$$(127) \quad 1 = \|R\| = \int_0^\infty |R(r)|^2 r dr$$

The **standard deviation of the nuclear charge distribution** in the core zone in natural units is:

$$(128) \quad s \equiv \frac{\bar{s}}{\rho_0}$$

Multiplying both sides of (95) with  $a_0^3/e$  and substituting via (101), (104), (105) and (128) yields the **nuclear charge distribution of the core zone** in natural units:

$$(129) \quad \sigma_{n,c}(r) = \frac{Q_c \lambda_n}{2\pi e L s^2} \exp\left(-\frac{r^2}{2s^2}\right)$$

### 3.11 Approximate Solution of the Radial Wave Function

The following **ansatz** will be used for approximating the **radial wave function**:

$$(130) \quad R(r) = f(r) \cdot \exp(-\zeta r), \text{ where}$$

$f(r)$  is assumed to be a polynomial and  $\zeta \in \mathbf{R}^+$  is a tunable scaling factor.

The radial Klein-Gordon equation (117) has a second solution, which is linear independent of the solution gained by ansatz (130). The second solution would be represented by the following ansatz:

$$(131) \quad R(r) = f(r) \cdot \exp(\zeta r), \text{ where } \zeta \in \mathbf{R}^+$$

However, this second solution and all linear combinations with it were incompatible with boundary condition (33). Therefore, this second solution ansatz will not be used.

The first derivative of the radial wave functions (130) reads:

$$(132) \quad R' = (f' - \zeta f) \cdot \exp(-\zeta r)$$

The second derivative of the radial wave functions is:

$$(133) \quad R'' = (f'' - 2\zeta f' + \zeta^2 f) \cdot \exp(-\zeta r)$$

**The value of  $\zeta$  can be determined by analyzing the asymptotic behavior of the wave function  $R$  at  $r \rightarrow \infty$ :**

The electromagnetic potential (and therefore the terms  $E_C$  and  $A_z$ ) become zero, when the radius approaches infinity.

Also, the terms proportional to  $1/r$  and  $1/r^2$  disappear at  $r \rightarrow \infty$ . The Klein-Gordon equation (117) then simplifies to:

$$(134) \quad -\lambda_n R'' + (P_z^2 - \alpha^2 E^2 - 2E)R = 0$$

Inserting (131) and (133) into (134) leads to:

$$(135) \quad \lambda_n (-f'' + 2\zeta f' - \zeta^2 f) + (P_z^2 - \alpha^2 E^2 - 2E)f = 0$$

Assuming function  $f$  can be approximated by a polynomial of finite degree, the function dominates over its derivatives at  $r \rightarrow \infty$  and therefore the **exponential scaling factor** is:

$$(136) \quad \zeta = \sqrt{\frac{1}{\lambda_n} (P_z^2 - \alpha^2 E^2 - 2E)}, \text{ where } E < 0$$



The case of  $E \geq 0$  does not apply here, because of simplification (21). Only the positive value of the square root is valid here, because of ansatz (130).

Solving (136) for the energy provides:

$$(137) \quad E = \frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2} + \frac{1}{\alpha^2} - \frac{1}{\alpha^2}$$

Only the positive value of the square root is valid here, because of boundary condition (37).

Equation (137) in conjunction with boundary condition (37) has interesting consequences:

$$(138) \quad \frac{|P_z|}{\sqrt{\lambda_n}} < \zeta < \sqrt{\frac{1}{\lambda_n} \left( P_z^2 + \frac{1}{\alpha^2} \right)}, \text{ i.e. the scaling factor } \zeta \text{ is limited by means of the axial canonical momentum at the low end and by means of the rest energy at the high end}$$

At the **non-relativistic limit** the **exponential scaling factor** computes as:

$$(139) \quad \zeta = \sqrt{\frac{1}{\lambda_n} (P_z^2 - 2E)}, \text{ thus } E = \frac{1}{2} (P_z^2 - \lambda_n \zeta^2)$$

Inserting (130), (132), (133) and (137) into radial Klein-Gordon equation (117) is leading to:

$$(140) \quad -\frac{\lambda_n}{2} f'' + \frac{\lambda_n}{2} \left( 2\zeta - \frac{1}{r} \right) f' + \left[ \frac{\lambda_n \zeta}{2r} + \frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left( \frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2} + \frac{1}{\alpha^2} - E_C \right)^2 + \frac{1}{2\alpha^2} - \frac{\lambda_n \zeta^2}{2} \right] f = 0$$

For the non-relativistic limit the Schrödinger equation (118) is leading to:

$$(141) \quad -\frac{\lambda_n}{2} f'' + \frac{\lambda_n}{2} \left( 2\zeta - \frac{1}{r} \right) f' + \left( \frac{\lambda_n \zeta}{2r} + \frac{\lambda_n m^2}{2r^2} - P_z A_z + \frac{A_z^2}{2} + E_C \right) f = 0$$

Solutions to differential equation (140) or (141) consist of eigenvalues of  $\zeta$  and eigenstates of polynomial  $f$ . These solutions can then be used to compute the eigenvalues of  $E$  and eigenstates of  $R$  of the radial Klein-Gordon equation (117) or Schrödinger equation (118).

**Function  $f(r)$**  can be **approximated by a polynomial** of  $r$  as following:

$$(142) \quad f(r) \approx \sum_{j=0}^J c_j r^{\beta+j} \text{ for } \beta \in \mathbf{N}_0 \text{ and } c_j \in \mathbf{R}$$

The summation runs over a number  $J+1$  of terms, depending on the desired accuracy of the approximation (in practice  $J$  needs to be up to 500 with 80-bit floating point numbers for “reasonable” accuracy). The (generally arbitrary) phase of the (generally complex) wave function  $R$  is chosen, such that the coefficients  $c_j$  become real numbers.

Generally, constants  $c_j$  and  $\zeta$  are depending on quantum numbers  $n$ ,  $m$  and  $l$ . For simplicity reasons, this dependency is not reflected in the respective indices of these constants.

The first derivative of (142) reads:

$$(143) \quad f'(r) \approx \sum_{j=0}^J (\beta + j) c_j r^{\beta+j-1}$$

The second derivative of (142) is:

$$(144) \quad f''(r) \approx \sum_{j=0}^J (\beta + j)(\beta + j - 1) c_j r^{\beta+j-2}$$

In equation (140) a number of terms can be approximated by a polynomial of degree  $P$ :

$$(145) \quad \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left( \frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2 + \frac{1}{\alpha^2}} - E_C \right)^2 + \frac{1}{2\alpha^2} - \frac{\lambda_n \zeta^2}{2} \approx \sum_{p=0}^P b_p r^p,$$

where  $P \leq J - 1$  and  $b_p \in \mathbf{R}$

At the non-relativistic limit (145) simplifies to:

$$(146) \quad -P_z A_z + \frac{A_z^2}{2} + E_C \approx \sum_{p=0}^P b_p r^p, \text{ where } P \leq J - 1 \text{ and } b_p \in \mathbf{R}$$

Approximations (145) and (146) probably have a limited convergence radius, no matter how large  $P$  is made and how the coefficients are chosen. However, for a given closed interval of radius values the approximations can be made arbitrarily precise by choosing  $P$  and the coefficients appropriately.

A suitable approximation can be found by first determining the range  $r_0 \leq r \leq r_p$  of relevant radius values reflecting the radial extent of the electron's wave function. For example, one can choose  $r_0$  and  $r_p$  in such a way, that the electron resides with 99.9% probability between these radii and, at the same time, the range is made as small as possible.

Based on this range, additional nodes  $r_1$  through  $r_{p-1}$  need to be determined between  $r_0$  and  $r_p$ . The nodes should be chosen, such that the approximation error is minimized (e.g. via Chebyshev nodes). These nodes can then be used e.g. by Newton polynomials for interpolation.

Inserting (142), (143), (144) and (145) into (140) and multiplying both sides with  $-2r^2/\lambda_n$  yields:

$$(147) \quad \sum_{j=0}^J (\beta + j)(\beta + j - 1) c_j r^{\beta+j} + (1 - 2\zeta r) \sum_{j=0}^J (\beta + j) c_j r^{\beta+j} - \left[ \zeta r + m^2 + \frac{2}{\lambda_n} \sum_{p=0}^P b_p r^{p+2} \right] \sum_{j=0}^J c_j r^{\beta+j} = 0$$

By neglecting the terms with potencies of  $r$  higher than  $\beta + J$  the result can be written as:

$$(148) \quad \sum_{j=0}^J \left\{ [(\beta + j)^2 - m^2] c_j + \zeta (1 - 2\beta - 2j) c_{j-1} - \frac{2}{\lambda_n} \sum_{p=0}^P b_p c_{j-p-2} \right\} r^{\beta+j} + \varepsilon_{cut} = 0,$$

where  $c_i = 0$  for  $i < 0$

and  $\varepsilon_{cut}$  is the cut-off error produced by neglecting potencies of  $r$  higher than  $\beta + J$

The cut-off error computes as:

$$(149) \quad \varepsilon_{cut} = \zeta (-1 - 2\beta - 2J) c_J r^{\beta+J+1} - \frac{2}{\lambda_n} \sum_{p=0}^P \sum_{j=J+1}^{J+p+1} b_p c_{j-p-2} r^{\beta+j}$$

The left hand side of equation (148) equals zero for all values of  $r$ . This can only be true, if the coefficients of  $r^{\beta+j}$  fulfill the following equation:

$$(150) \quad [(\beta + j)^2 - m^2] c_j + \zeta (1 - 2\beta - 2j) c_{j-1} - \frac{2}{\lambda_n} \sum_{p=0}^P b_p c_{j-p-2} = 0$$

Analyzing the case  $j = 0$  gives:

$$(151) \quad \beta^2 - m^2 = 0, \text{ therefore } \beta = |m|$$

Inserting (151) into (150) yields the **iterative formula for computing the coefficients** from the value of  $c_0$ :

$$(152) \quad c_j = \frac{1}{(2|m|j + j^2)} \left\{ \zeta(2|m| + 2j - 1)c_{j-1} + \frac{2}{\lambda_n} \sum_{p=0}^P b_p c_{j-p-2} \right\}, \text{ where } c_i = 0 \text{ for } i < 0$$

Note, that the coefficients  $c_j$  are all proportional to each other. Formula (152) stays the same at the non-relativistic limit.

Equation (149) puts additional requirements on the coefficients  $c_{J-P}$  through  $c_J$ , which contradict the requirements of equation (152). Therefore, the polynomial approximation of the radial wave function with finite  $J$  cannot be made precise.

**The approximation error becomes minimal, when the last coefficient  $c_J$  is zero, which is the case only for the eigenvalues of  $\xi$ . Therefore, this defines a method for determining the eigenvalues.**

Alternatively one could determine the eigenvalues by using the original Klein-Gordon equation (117) in conjunction with (137) as a measure of error:

$$(153) \quad \delta(r, \zeta) = -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[ \frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} - \frac{\alpha^2}{2} \left( \frac{1}{\alpha} \sqrt{P_z^2 - \lambda_n \zeta^2 + \frac{1}{\alpha^2}} - E_C \right)^2 + \frac{1}{2\alpha^2} \right] R \\ \approx -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left\{ \frac{\lambda_n m^2}{2r^2} + \sum_{p=0}^P b_p r^p + \frac{\lambda_n \zeta^2}{2} \right\} R$$

At the non-relativistic limit one would use the Schrödinger equation (118) in conjunction with (139) as a measure of error:

$$(154) \quad \delta(r, \zeta) = -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left[ \frac{\lambda_n m^2}{2r^2} - P_z A_z + \frac{A_z^2}{2} + E_C + \frac{\lambda_n \zeta^2}{2} \right] R \\ \approx -\frac{\lambda_n}{2} R'' - \frac{\lambda_n}{2r} R' + \left\{ \frac{\lambda_n m^2}{2r^2} + \sum_{p=0}^P b_p r^p + \frac{\lambda_n \zeta^2}{2} \right\} R$$

The function  $\delta(r, E)$  is approaching zero for all values of  $r$  only at the energy eigenvalues  $E$  or  $\zeta$ .

**The value of coefficient  $c_0$  can be determined from  $\zeta$  by normalization of the wave function  $R$ .**

Combining (130), (142) and (151) leads to:

$$(155) \quad R \approx \sum_{j=0}^J c_j r^{|m|+j} \exp(-\zeta r)$$

The normalization condition (127) requires:

$$(156) \quad 1 = \|R\|^2 = \int_0^\infty |R(r)|^2 r dr = \int_0^\infty \left( \sum_{j=0}^J c_j r^{|m|+j} \right)^2 \exp(-2\zeta r) r dr$$

That means, one has to scale all  $c_j$  proportionally, such that (156) yields the value 1.

### 3.12 Total Binding Energy

One could naively assume, that the total binding energy  $E_B$  of a CP is the sum of the energies of all electrons plus the nuclear self-repulsion energy:

$$(157) \quad E_B \neq E_n + \sum_{i=1}^N E_i, \text{ where } E_i \text{ is the energy eigenvalue of electron number } i$$

Unfortunately, this approach would count the electron-electron interaction energies  $E_{C,e}$  and  $E_M$  twice, because the electrons are interacting with themselves.

Instead, the following method will be used for computing the total binding energy: In the first step, the energy eigenvalues  $E_i$  will be decomposed via the radial Klein-Gordon equation (117) to the expectation values of each energy term therein. Secondly, the expectation values containing the interaction energies  $E_{C,e}$  in a non-quadratic form will be divided by two and the expectation values containing these interaction energies in a quadratic form will be divided by four. Finally, with this adjustment, the corrected energies  $\tilde{E}_i$  will be reassembled via the same Klein-Gordon equation.

The **total binding energy of a CP** (in units of  $E_n/\alpha$ ) then computes as:

$$(158) \quad E_B = E_n + \sum_{i=1}^N \tilde{E}_i, \text{ where } \tilde{E}_i \text{ is the corrected energy of electron number } i$$

Multiplying (117) by  $R^* r$ , using (109) and expanding the squared parenthesis provides:

$$(159) \quad \frac{\lambda_n R^*}{2} (-rR'' - R') + |R|^2 r \left[ \frac{\lambda_n m^2}{2r^2} + \frac{p_z^2}{2} - \frac{\alpha^2 E_i^2}{2} - \frac{\alpha^2 E_{C,n}^2}{2} - \alpha^2 E_{C,n} E_{C,e} - \frac{\alpha^2 E_{C,e}^2}{2} + \alpha^2 E_i E_{C,n} + \alpha^2 E_i E_{C,e} - E_i + E_{C,n} + E_{C,e} \right] = 0$$

Integrating (159) over all radius values and using (127) transforms it to an equation of expectation values, thereby decomposing  $E_i$ :

$$(160) \quad E_i = \langle E_{C,n} \rangle + \langle E_{C,e} \rangle - \frac{1}{\alpha^2} \pm sqrt \left[ \left( -\langle E_{C,n} \rangle - \langle E_{C,e} \rangle + \frac{1}{\alpha^2} \right)^2 - \langle E_{C,n}^2 \rangle - 2\langle E_{C,n} E_{C,e} \rangle - \langle E_{C,e}^2 \rangle + \frac{2}{\alpha^2} \left( \langle E_r \rangle + \langle E_\varphi \rangle + \frac{1}{2} \langle p_z^2 \rangle + \langle E_{C,n} \rangle + \langle E_{C,e} \rangle \right) \right], \text{ where}$$

$$(161) \quad \langle E_r \rangle = \frac{\lambda_n}{2} \left\langle \frac{1}{R} \left( -R'' - \frac{R'}{r} \right) \right\rangle = \frac{\lambda_n}{2} \int_0^\infty R^* (-rR'' - R') dr \text{ is the term related to the radial kinetic energy,}$$

$$(162) \quad \langle E_\varphi \rangle = \frac{\lambda_n m^2}{2} \left\langle \frac{1}{r^2} \right\rangle = \frac{\lambda_n m^2}{2} \int_0^\infty \frac{|R|^2}{r} dr \text{ is the term related to the azimuthal kinetic energy,}$$

$$(163) \quad \frac{1}{2} \langle p_z^2 \rangle = \frac{1}{2} \langle (P_z - A_z)^2 \rangle = \frac{1}{2} \int_0^\infty \left( \frac{2\pi}{L} l - A_z \right)^2 |R|^2 r dr \text{ is the term related to the axial kinetic energy,}$$

$$(164) \quad \langle E_{C,n}^2 \rangle = \int_0^\infty E_{C,n}^2 |R|^2 r dr \text{ is the square of the Coulomb energy from electron-nucleus interaction,}$$

$$(165) \quad \langle E_{C,e}^2 \rangle = \int_0^\infty E_{C,e}^2 |R|^2 r dr \quad \text{is the square of the Coulomb energy from electron-electron interaction,}$$

and so on.

Dividing the expectation values in (160) containing the interaction energy  $E_{C,e}$  in a non-quadratic form by two and dividing the expectation values containing this interaction energy in a quadratic form by four yields the **corrected energies**:

$$(166) \quad \tilde{E}_i = \langle E_{C,n} \rangle + \frac{\langle E_{C,e} \rangle}{2} - \frac{1}{\alpha^2} \pm \text{sqr}t \left[ \left( -\langle E_{C,n} \rangle - \frac{\langle E_{C,e} \rangle}{2} + \frac{1}{\alpha^2} \right)^2 - \langle E_{C,n}^2 \rangle - \langle E_{C,n} E_{C,e} \rangle - \frac{\langle E_{C,e}^2 \rangle}{4} \right. \\ \left. + \frac{2}{\alpha^2} \left( \langle E_r \rangle + \langle E_\phi \rangle + \frac{1}{2} \langle p_z^2 \rangle + \langle E_{C,n} \rangle + \frac{1}{2} \langle E_{C,e} \rangle \right) \right]$$

The sign of the square root in (160) and (166) has to be negative if  $\langle E_{C,n} \rangle + \langle E_{C,e} \rangle > 0$  or positive otherwise.

At the **non-relativistic limit** the computation of the total binding energy is simpler. Multiplying (118) by  $R^* r$  and expanding the squared parenthesis provides:

$$(167) \quad \frac{\lambda_n R^*}{2} (-rR'' - R') + r|R|^2 \left[ \frac{\lambda_n m^2}{2r^2} + \frac{1}{2} p_z^2 + E_C - E_i \right] = 0$$

Integrating (167) over all radius values transforms it to an equation of expectation values, thereby decomposing  $E_i$ :

$$(168) \quad E_i = \langle E_r \rangle + \langle E_\phi \rangle + \frac{1}{2} \langle p_z^2 \rangle + \langle E_{C,n} \rangle + \langle E_{C,e} \rangle$$

Hence the corrected energies at the non-relativistic limit become:

$$(169) \quad \tilde{E}_i = \langle E_r \rangle + \langle E_\phi \rangle + \frac{1}{2} \langle p_z^2 \rangle + \langle E_{C,n} \rangle + \frac{1}{2} \langle E_{C,e} \rangle$$

### 3.13 Grouping, Orbital Occupation, Self-Consistent Field Iterations

The electron configuration of a CP consists of many orbitals, which are characterized by the quantum numbers  $n$ ,  $m$  and  $l$ . According to the Pauli exclusion principle each orbital can only be occupied by a maximum of two electrons (one with spin up and one with spin down).

There are too many electrons in a CP for computing all occupied orbitals individually. Instead, ranges of orbitals with contiguous values for  $l$  are grouped together. Within a group all orbitals have the same quantum numbers  $n$  and  $m$ . These orbitals of such groups differ in quantum number  $l$ . The arithmetic mean of the quantum numbers  $l$  represents the group during computation.

The most simple approach is to let each group contain the same number of orbitals. On one hand the groups should be small enough to achieve a fine spacing in the electron energies (for accuracy). On the other hand the groups need to be coarse enough, such that computation time becomes affordable.

Equations (120) and (121) are computed by letting the summation run over the occupied number of groups. Each summand is multiplied by the number of electrons it represents.

For ground state computations the occupation should start with the lowest energy. It should progress to groups with successively higher energy until the targeted number of electrons “found their orbital”.

Equations (120), (121) and (152), as well as the occupation process are depending on each other in a circular manner. Thus they can be computed only **iteratively** until reaching **self-consistency between eigenstates, potential and occupation**.

According to simplification (19) orbitals with eigenvalue energies below  $-m_e c^2$  are “forbidden” to occupy. Care must be taken and appropriate numerical damping must be applied, such that fluctuations of orbitals between “forbidden” and “allowed” are not hindering convergence of the SCF algorithm.

Within each of these SCF-iterations (self-consistent field iterations) there is a need for sub-iterations:

According to (121) and (124) the axial current density  $J_z$  and the vector potential  $A_z$  are mutually depending on each other. Sub-iterations are required for making these quantities consistent with each other, while leaving the eigenstates unchanged.